

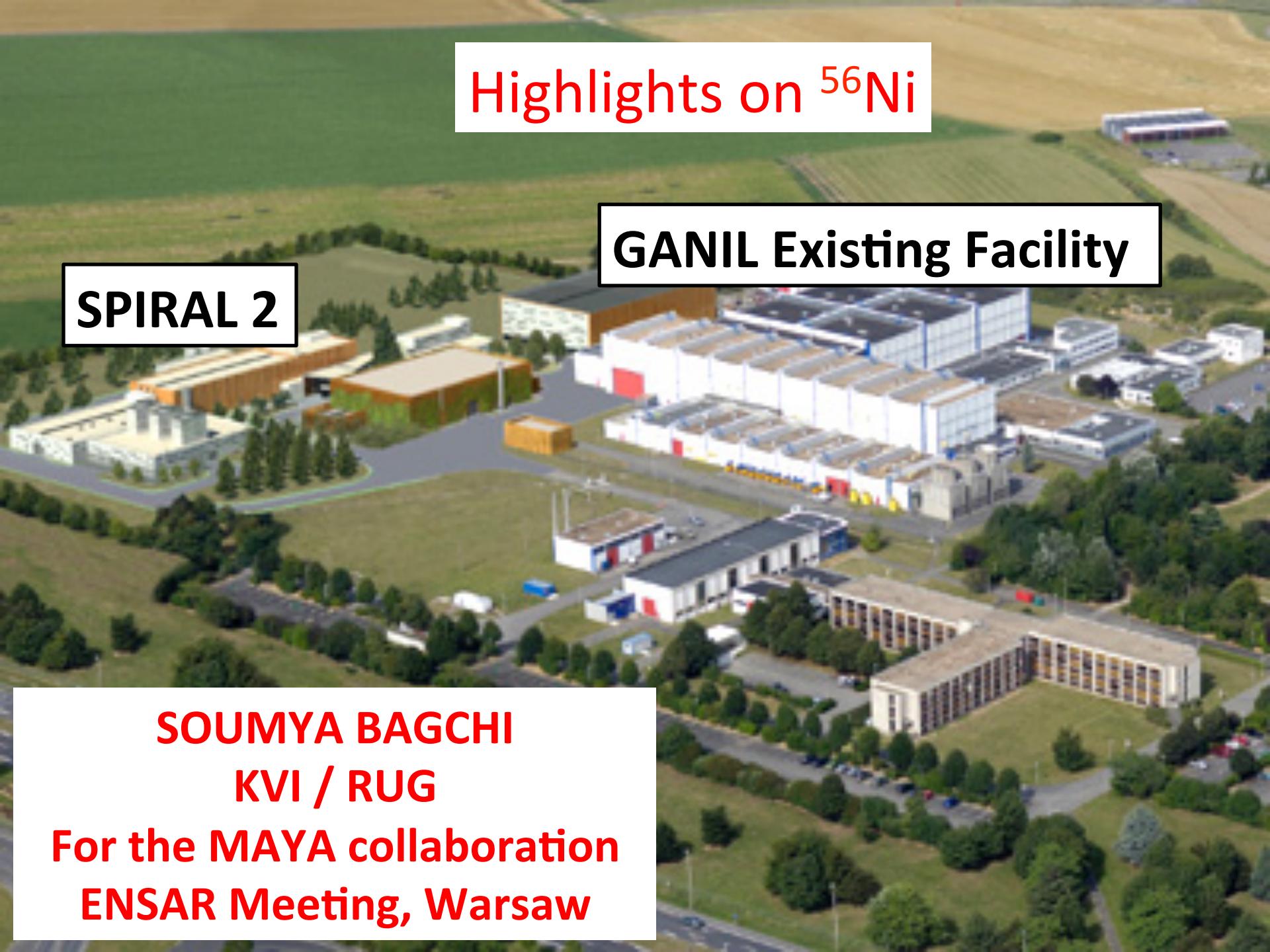
Highlights on ^{56}Ni

GANIL Existing Facility

SPIRAL 2

**SOUMYA BAGCHI
KVI / RUG**

**For the MAYA collaboration
ENSAR Meeting, Warsaw**



Outline:

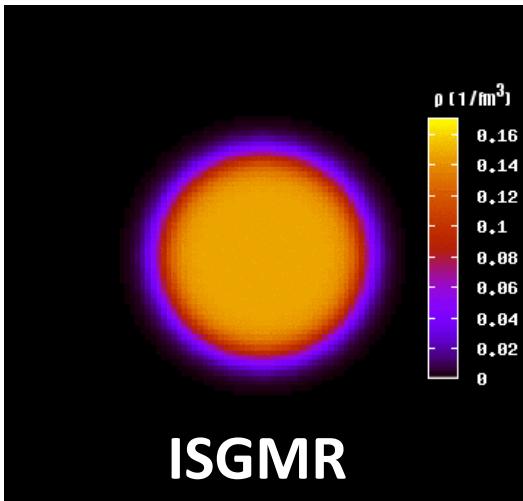
- Introduction to Giant Resonances
- Importance of compression modes in nuclei
- Experimental Setup (MAYA)
- Results
- Hint for alpha clustering (H. Akimune)
- Summary and outlook

Giant Resonances:

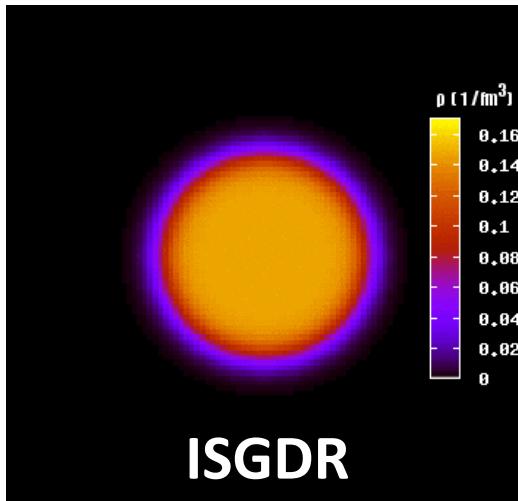
Collective oscillations of all neutrons and all protons in a nucleus

- In phase (Iso-Scalar)

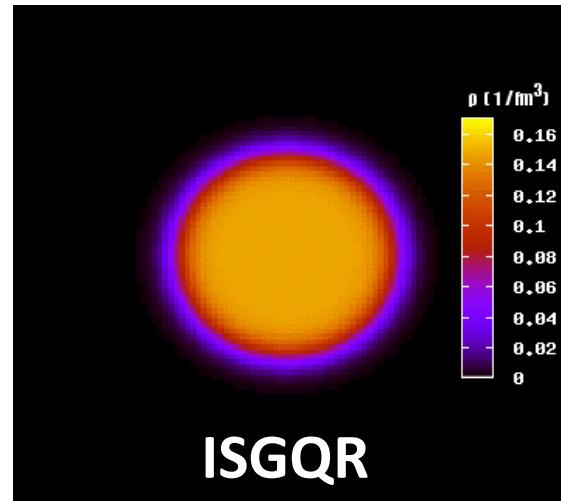
M. Itoh



Monopole, $L = 0$

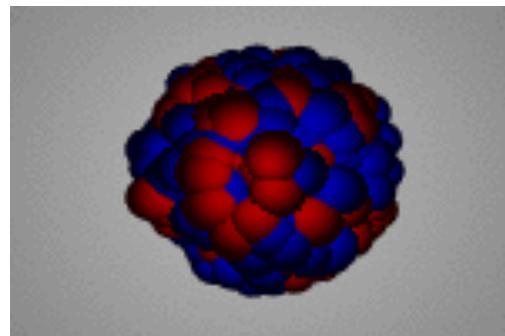


Dipole, $L = 1$



Quadrupole, $L = 2$

- Out of phase (Iso-Vector)



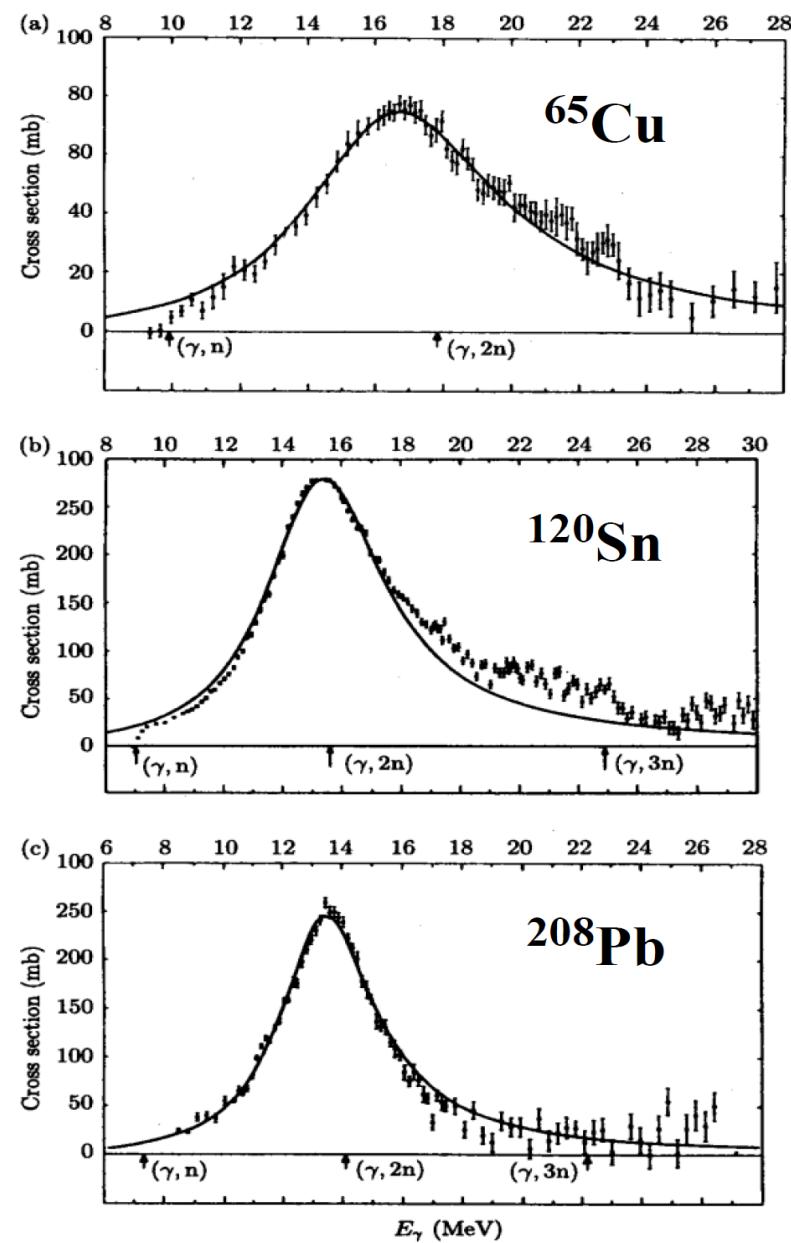
Macroscopic & Microscopic Structure

Giant resonances (collective modes)

- Width and Excitation energy
- Coherent superposition of 1p-1h excitation

(Figure depicts IVGDR, obtained by gamma absorption in these nuclei)

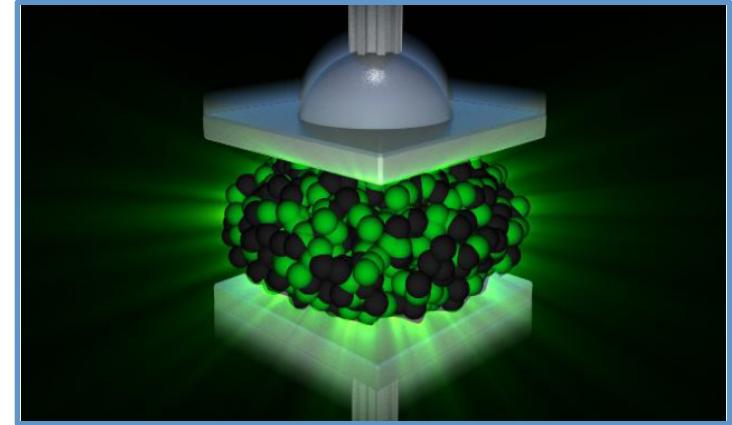
Berman and Fultz, Rev. Mod. Phys. 47 (1975)



Nuclear Incompressibility

Incompressibility → Measure of the resistance of matter to uniform compression

$$K = -V \frac{\partial P}{\partial V}$$



Adapted from M. N. Harakeh

ISGMR ($T=0, L=0$)



$$E_{ISGMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

ISGDR ($T=0, L=1$)



$$E_{ISGDR} = \hbar \sqrt{\frac{\frac{7}{3} K_A + \frac{27}{25} \varepsilon_F}{m \langle r^2 \rangle}}$$

Nuclear Incompressibility

$$K_A = K_{vol} + K_{surf} A^{-1/3} + K_\tau ((N-Z)/A)^2 + K_{Coul} Z^2 A^{-4/3}$$

EoS is important for studying:

- Core collapse and explosion of supernovae
- Formation of neutron stars
- Collisions of heavy ions

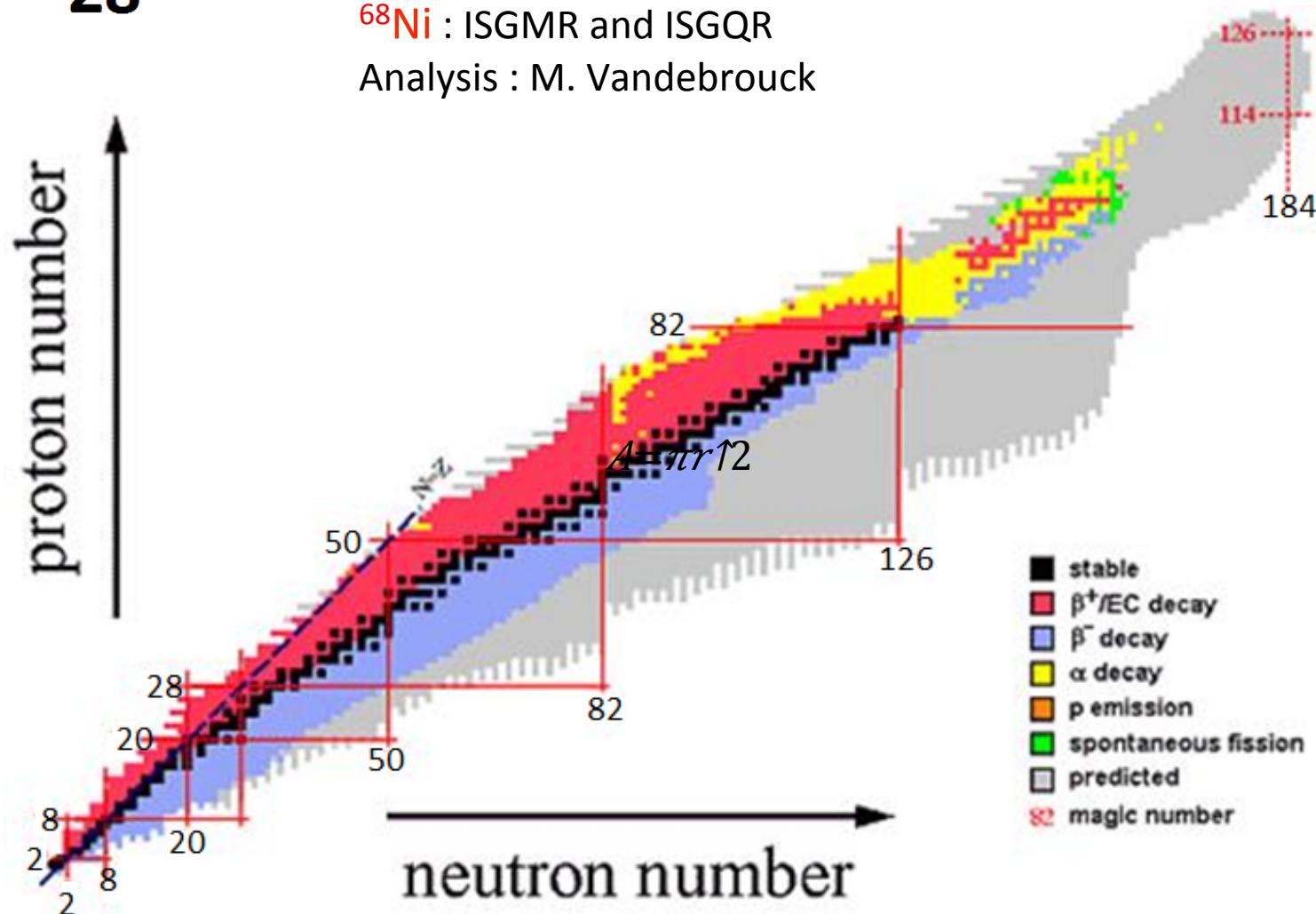
Why Ni ?

- Incompressibility value obtained for
Pb isotopes 240 ± 10 MeV
Sn and Cd isotopes $210 - 215$ MeV
- Question : Why Sn and Cd are softer than Pb?
- Need of study of incompressibility for a series of isotopes of a nucleus

$^{56}_{28}\text{Ni}_{28}$

^{56}Ni : ISGMR and ISGQR
C. Monrozeau et al., PRL 100, 042501 (2008)

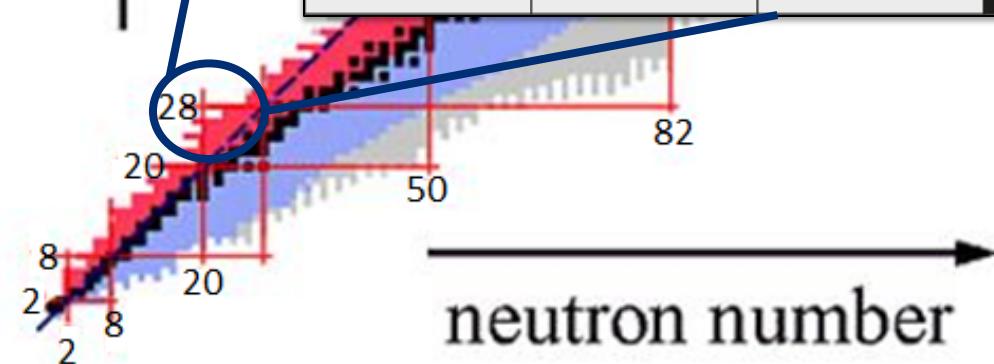
^{68}Ni : ISGMR and ISGQR
Analysis : M. Vandebrouck



$$K_A = K_{vol} + K_{surf} A^{-1/3} + K_\tau ((N-Z)/A)^2 + K_{Coul} Z^2 A^{-4/3}$$

$^{56}_{28}\text{Ni}_{28}$

proton number

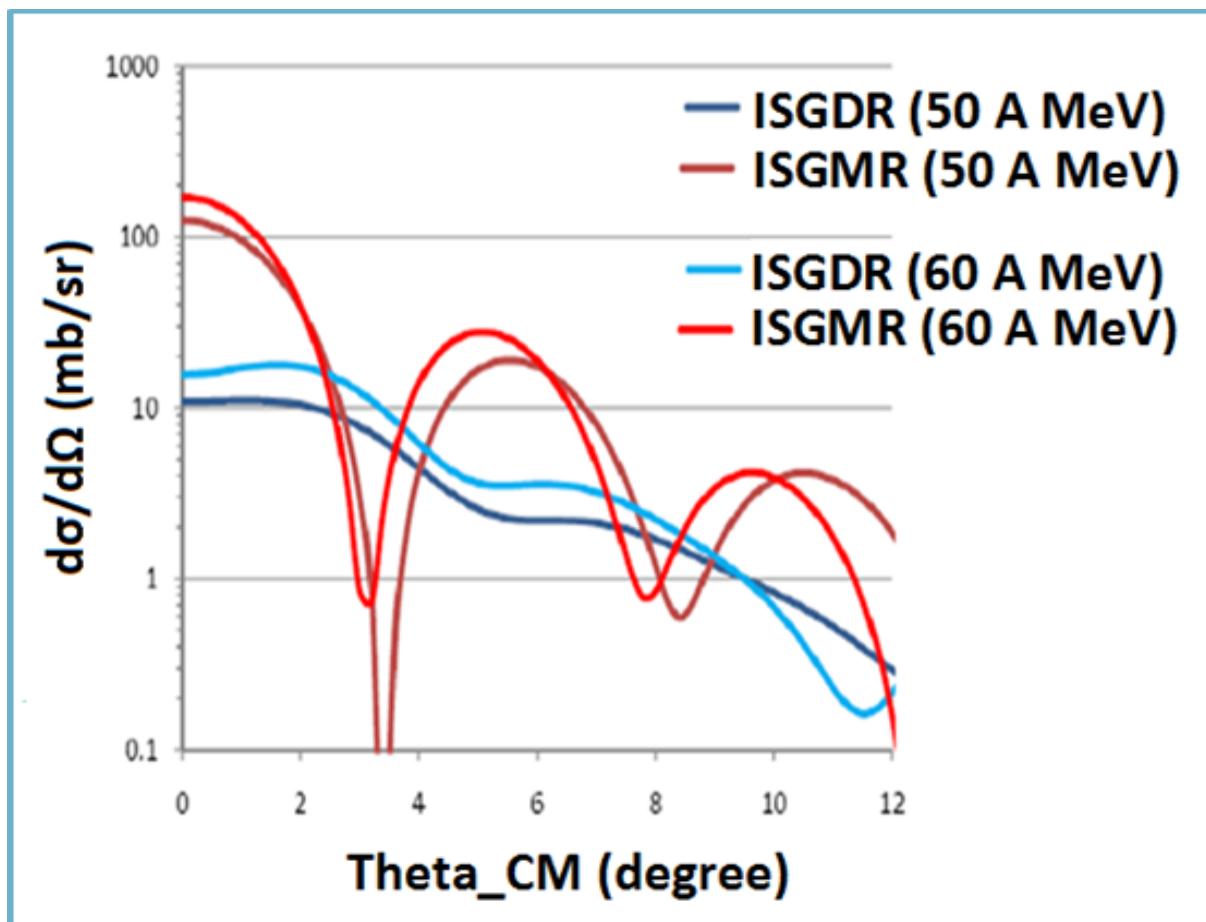


$^{53}_{28}\text{Ni}$ 45 ms	$^{54}_{28}\text{Ni}$ 104 ms	$^{55}_{28}\text{Ni}$ 205 ms	n	$^{56}_{28}\text{Ni}$ 6.08 d	$^{57}_{28}\text{Ni}$ 35.6 h	$^{58}_{28}\text{Ni}$ stable
$^{52}_{26}\text{Co}$ 115 ms	$^{53}_{26}\text{Co}$ 242 ms	$^{54}_{26}\text{Co}$ 193 ms	p	$^{55}_{26}\text{Co}$ 17.53 h	$^{56}_{26}\text{Co}$ 77 d	
$^{51}_{26}\text{Fe}$ 305 ms	$^{52}_{26}\text{Fe}$ 8.28 h	$^{53}_{26}\text{Fe}$ 8.51 m	α	$^{54}_{26}\text{Fe}$ stable		

- β^- decay
- α decay
- p emission
- spontaneous fission
- predicted
- magic number

Angular Distribution of ISGMR and ISGDR

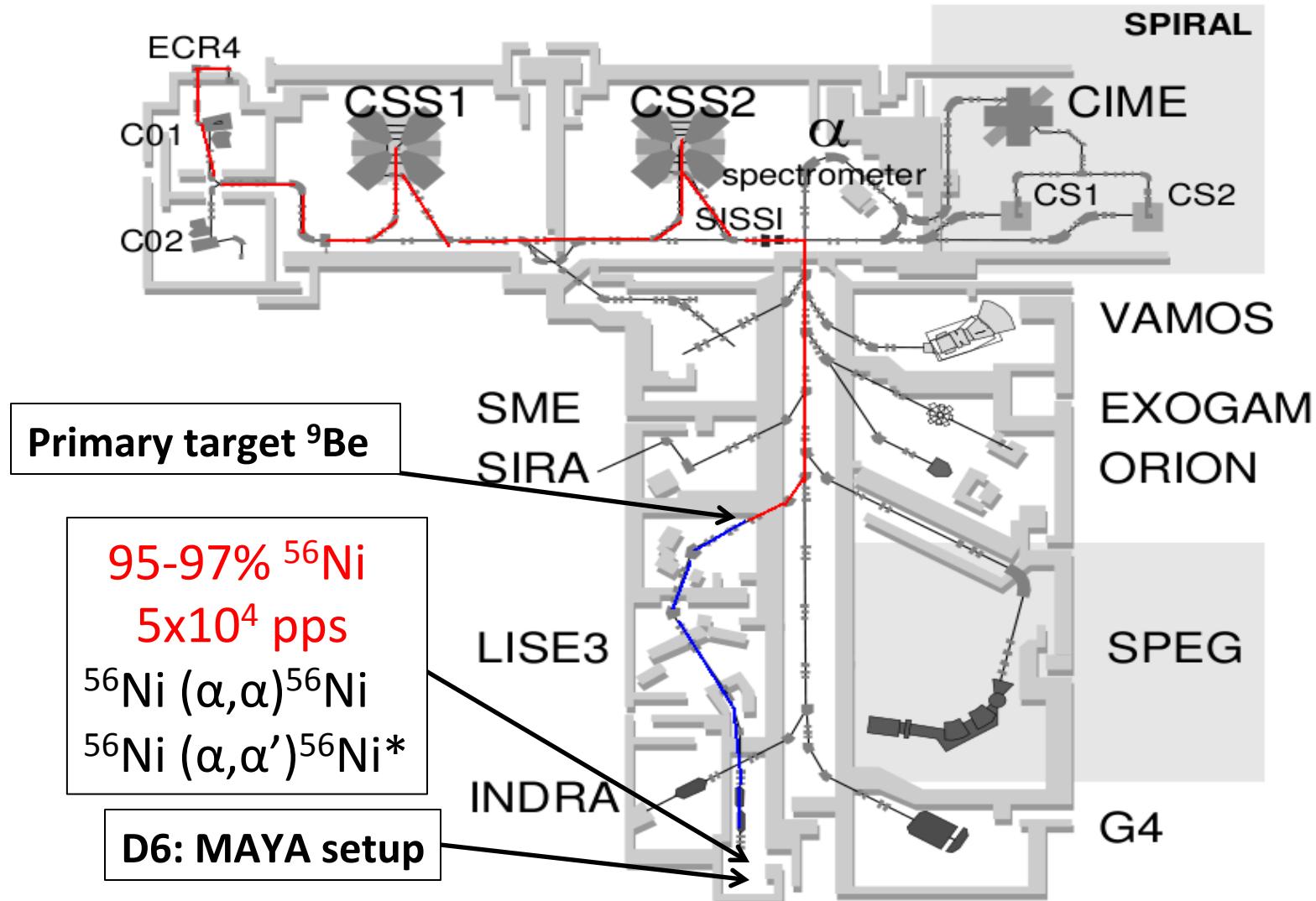
Differential cross-section in DWBA calculation for ^{56}Ni :



Primary Beam : ^{58}Ni at 75 MeV/u

Fragmentation method : Target ^9Be (thickness 500 μm)

Secondary Beam : ^{56}Ni at 50 MeV/u



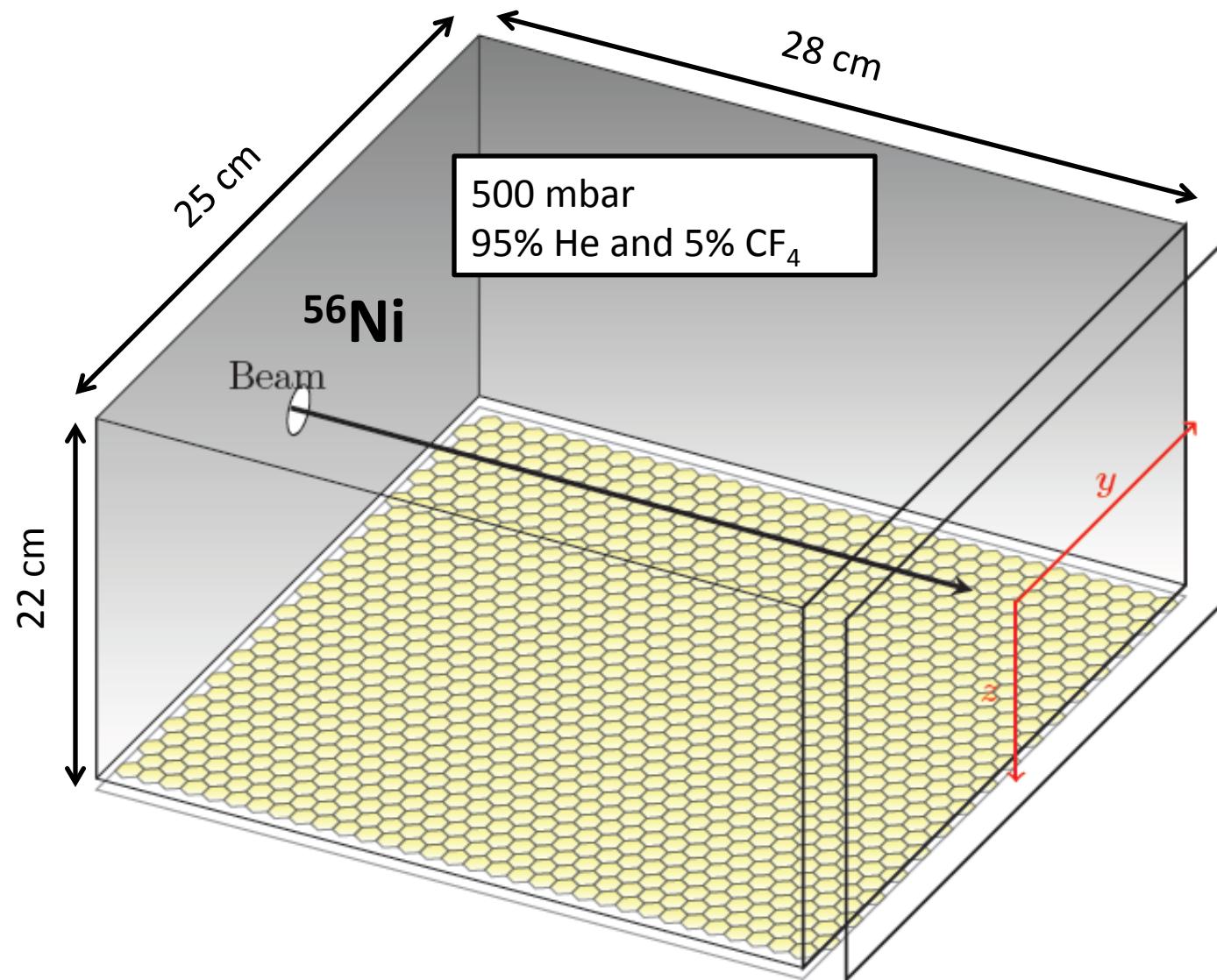
Challenges with exotic beams:

- Intensity of exotic beams is very low
- To get reasonable yields thick target is needed
- Very low energy (\sim sub MeV) recoil particle will not come out of the target

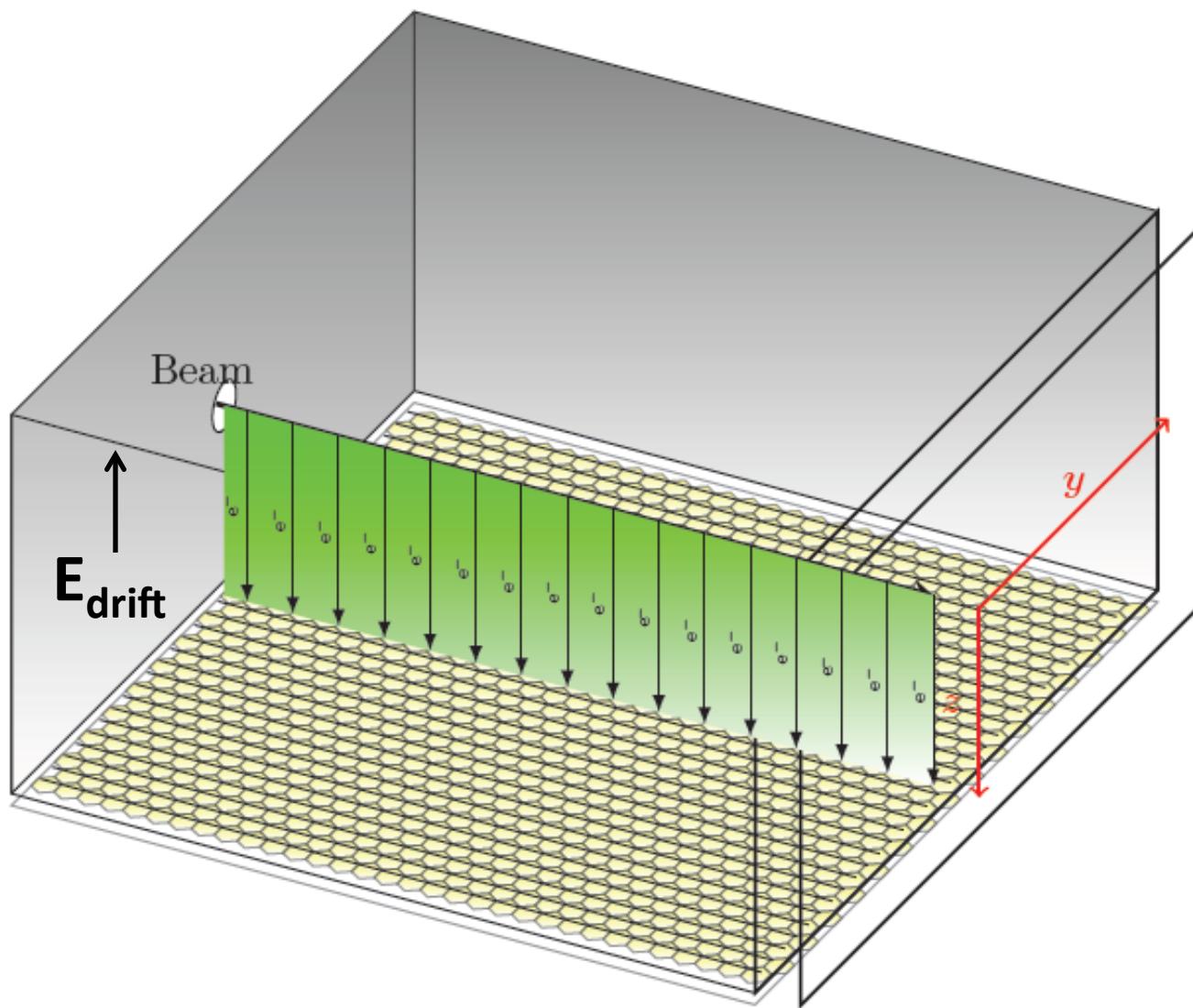
Active target: Detection takes place at every point of the target

- Good angular coverage
- Effective target thickness can be increased without much loss of resolution
- Detection of very low energy recoil particle is also possible

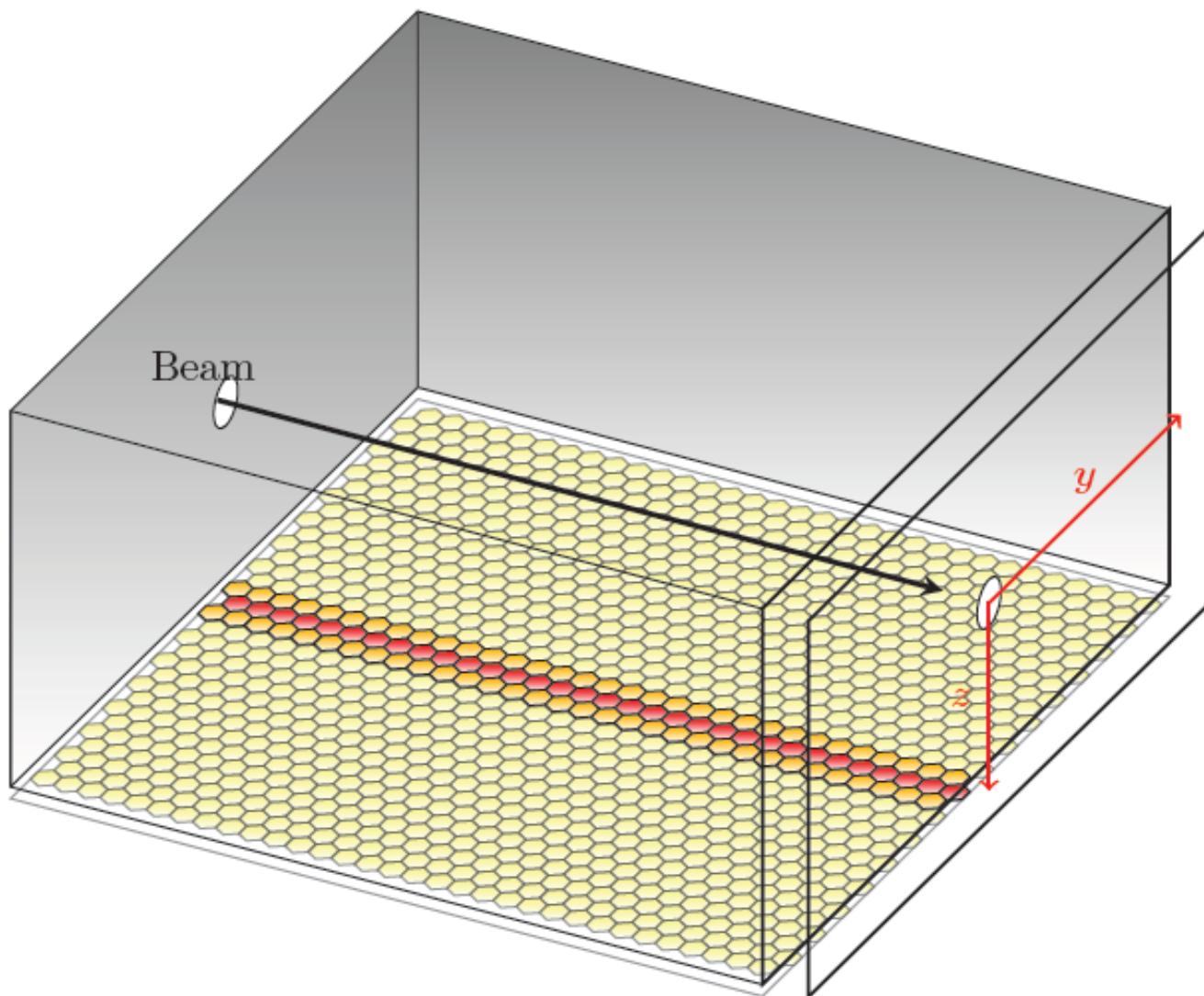
MAYA setup



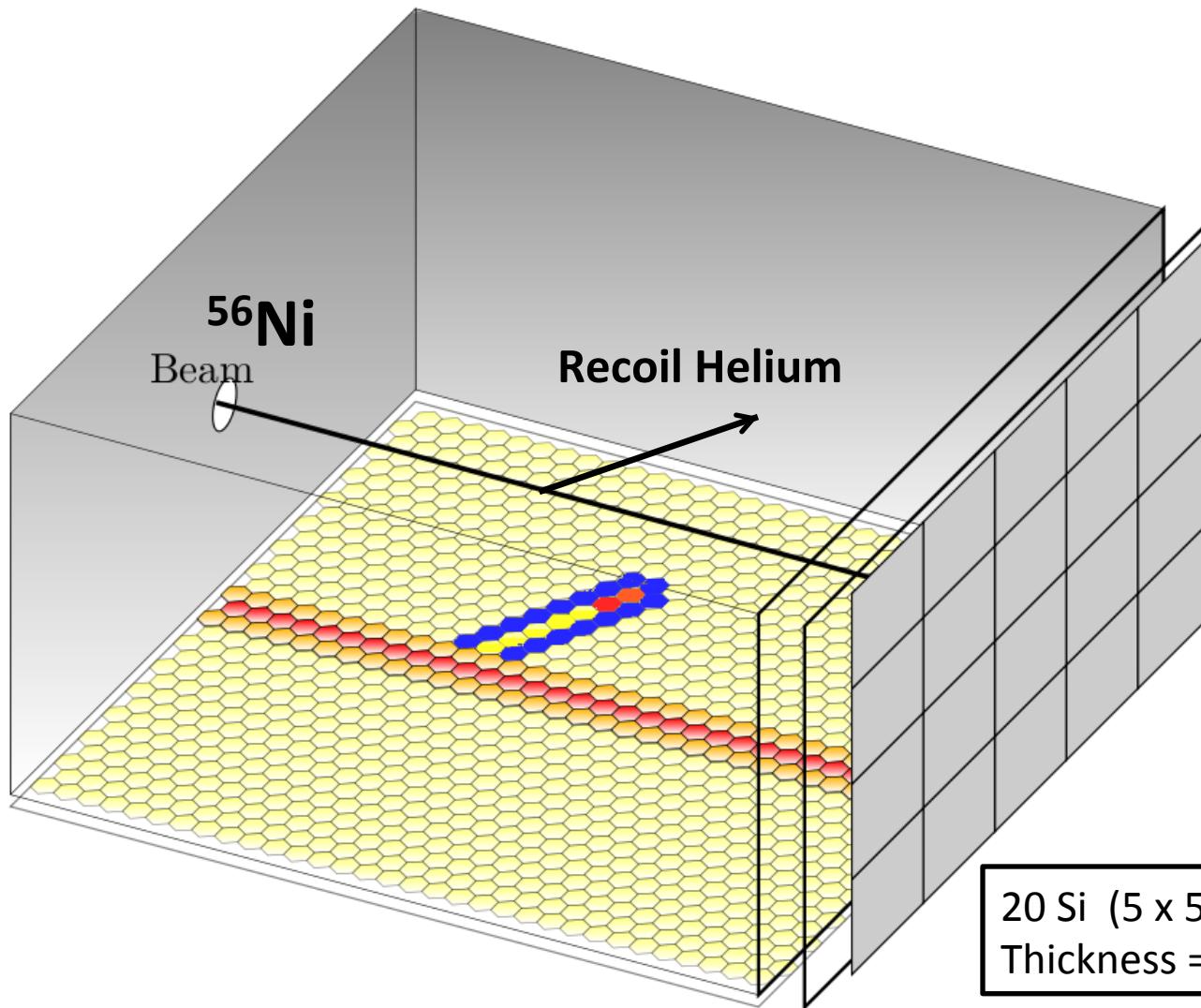
MAYA setup



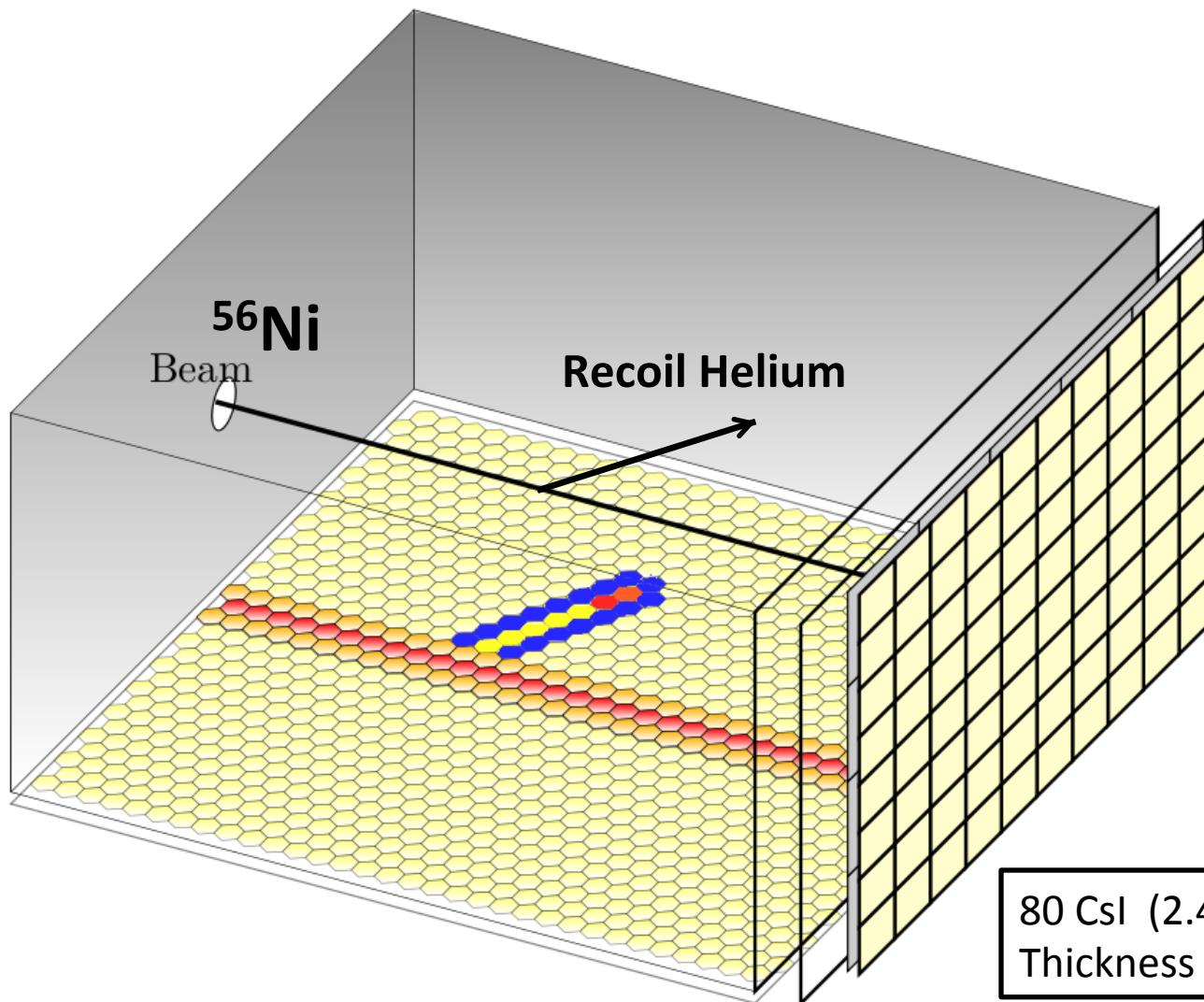
MAYA setup



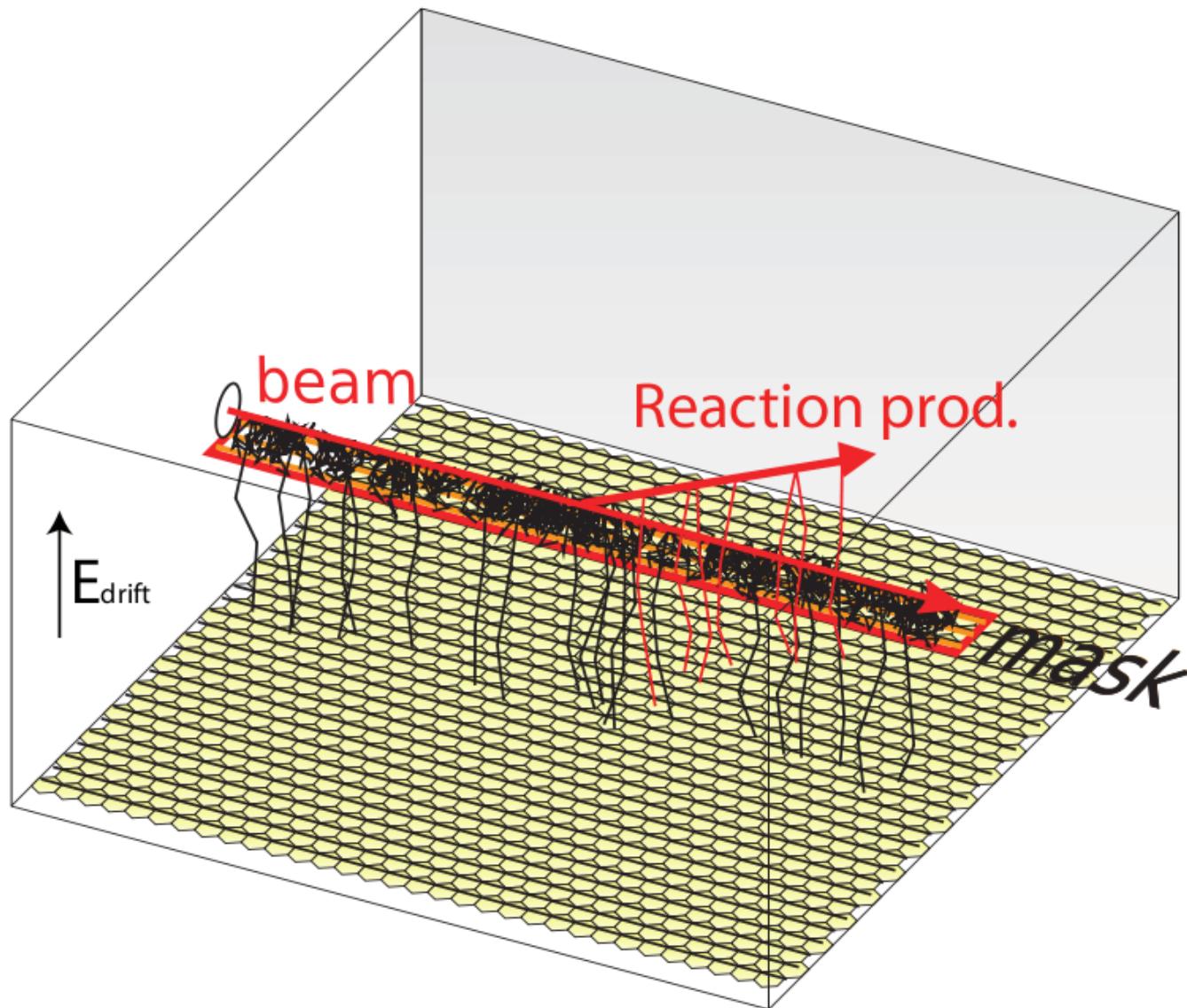
Si/CsI Telescope in MAYA

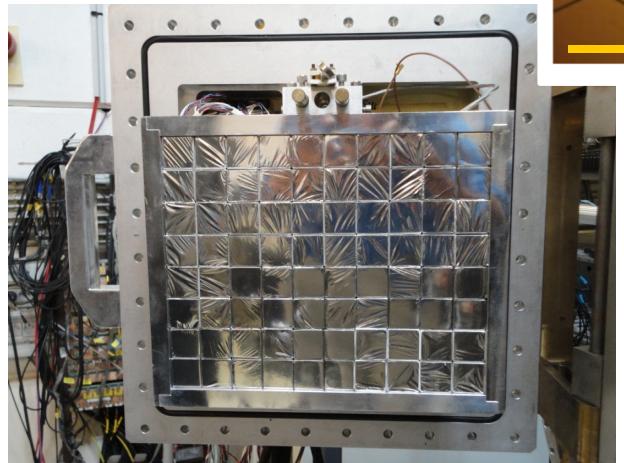
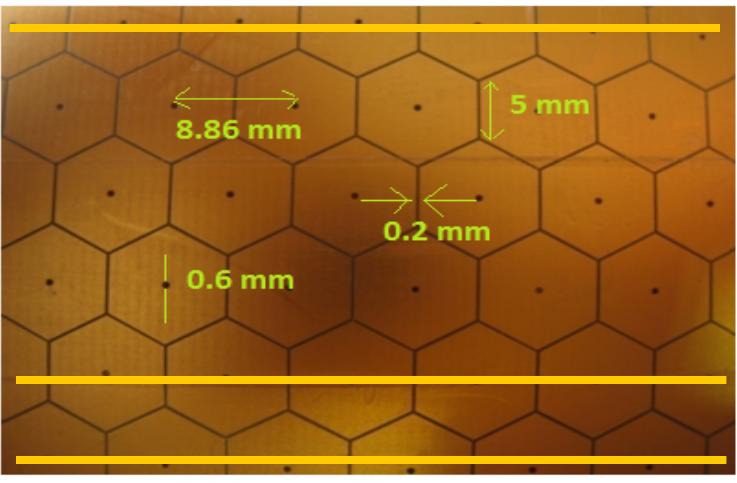
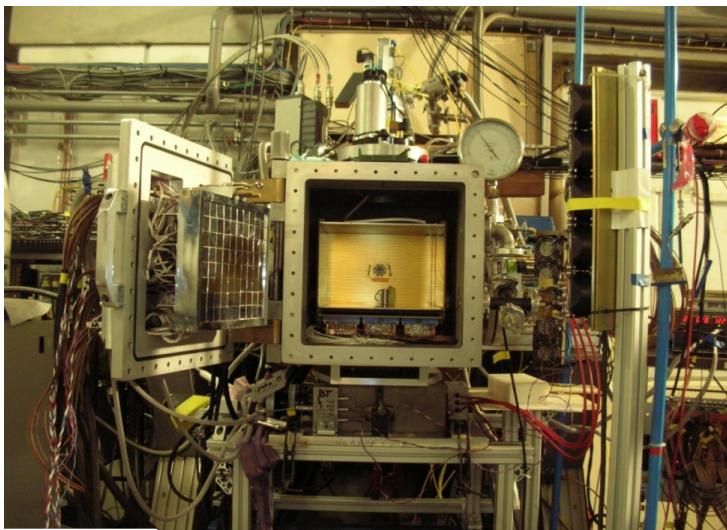
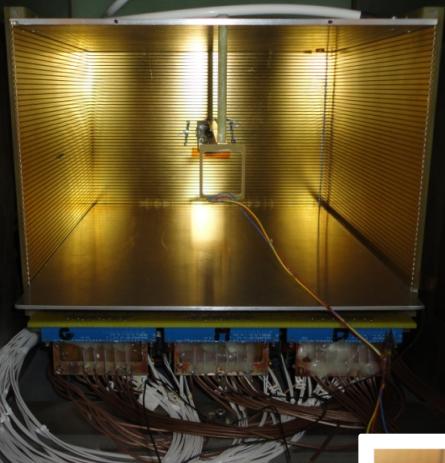


Si/CsI Telescope in MAYA

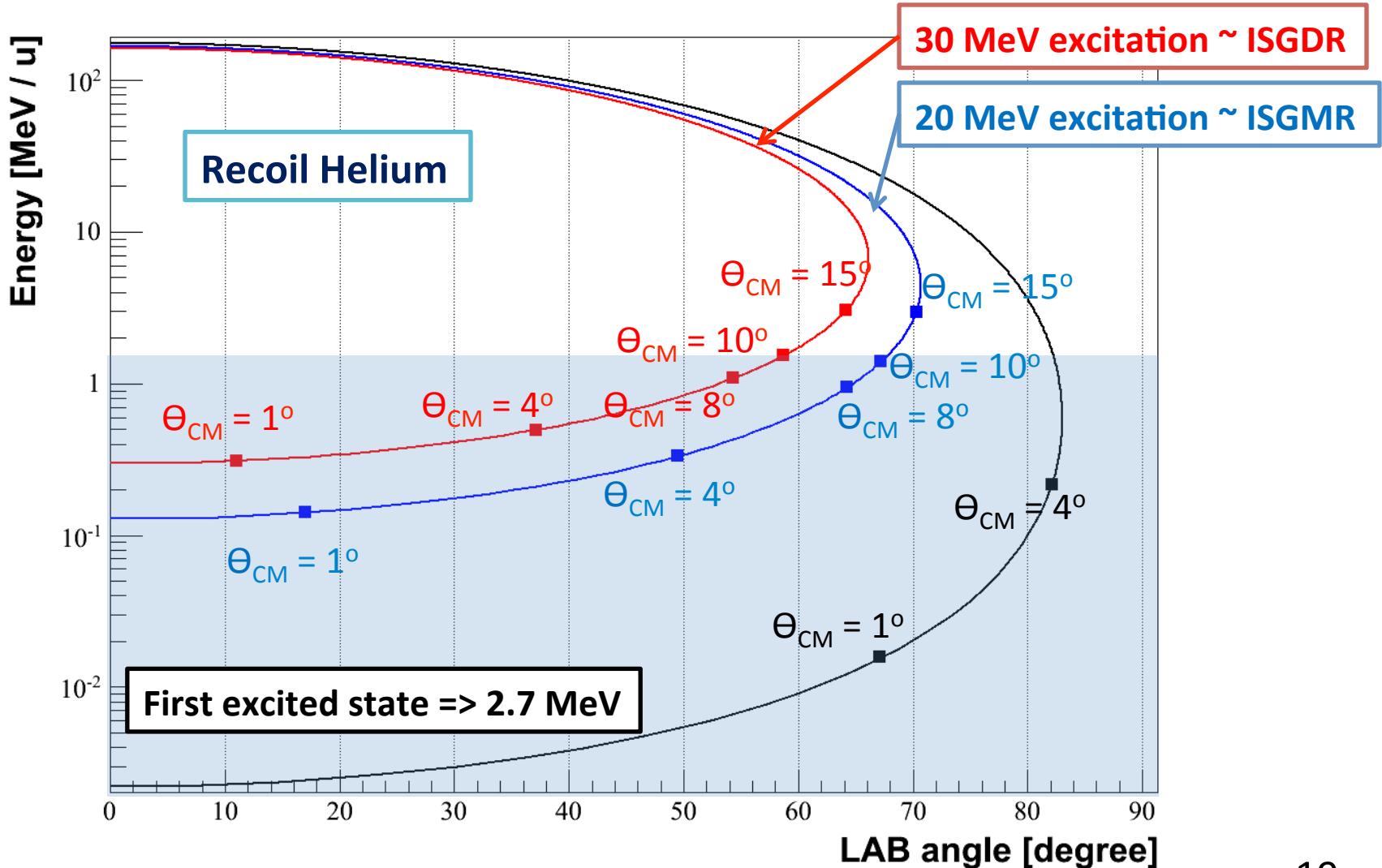


Mask in MAYA

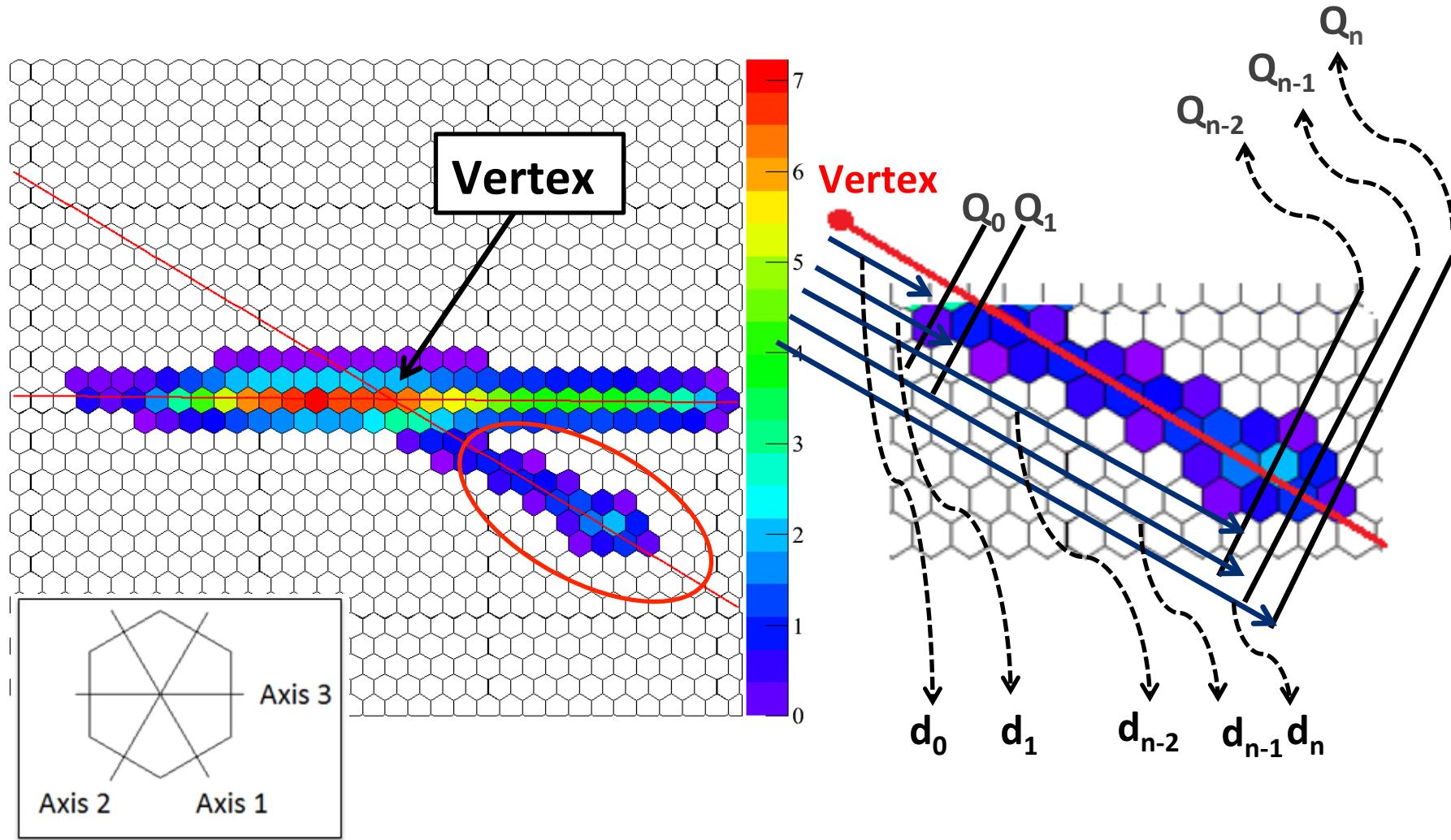




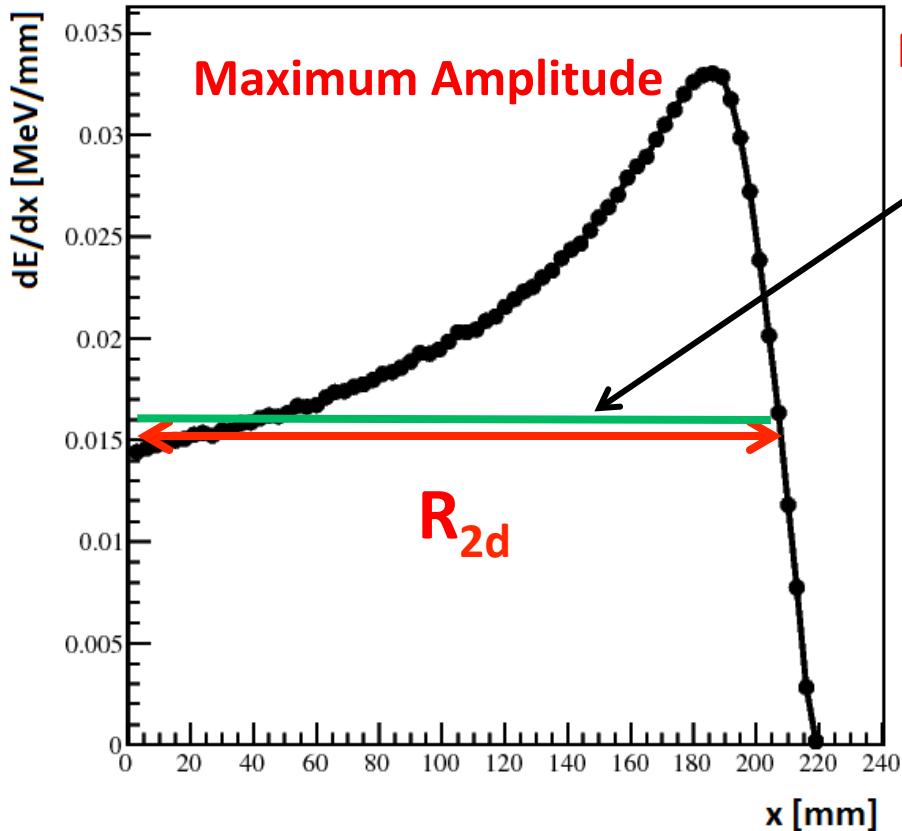
Kinematics curve for $^{56}\text{Ni} (\alpha, \alpha')$ $^{56}\text{Ni}^*$ with ^{56}Ni at 50 MeV / u



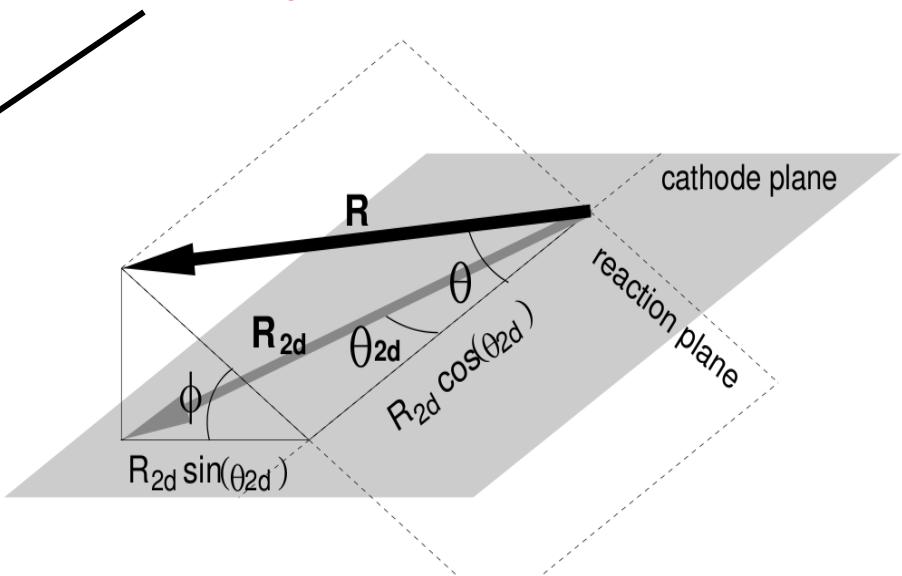
Kinematics reconstruction in MAYA



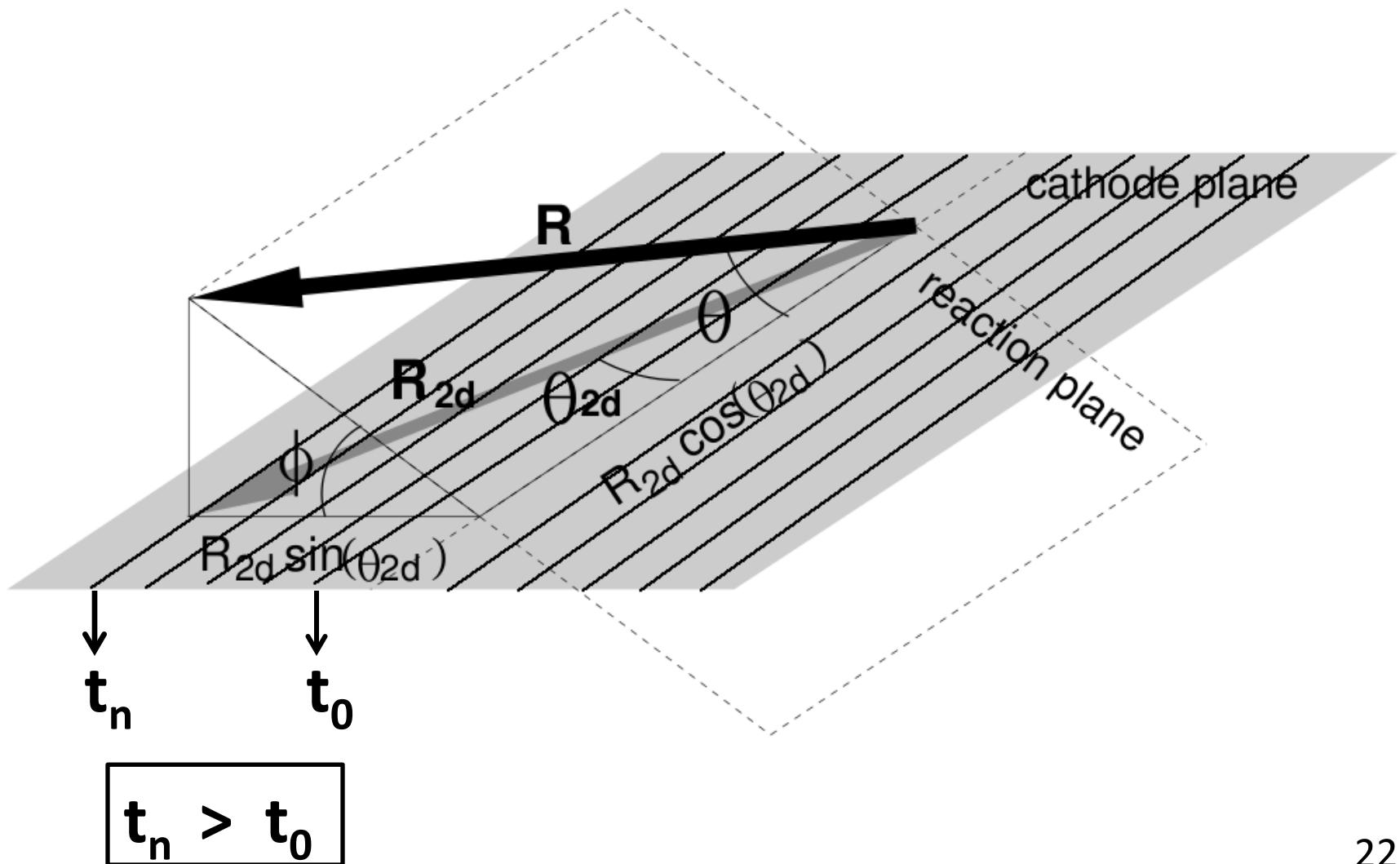
Kinematics reconstruction in MAYA



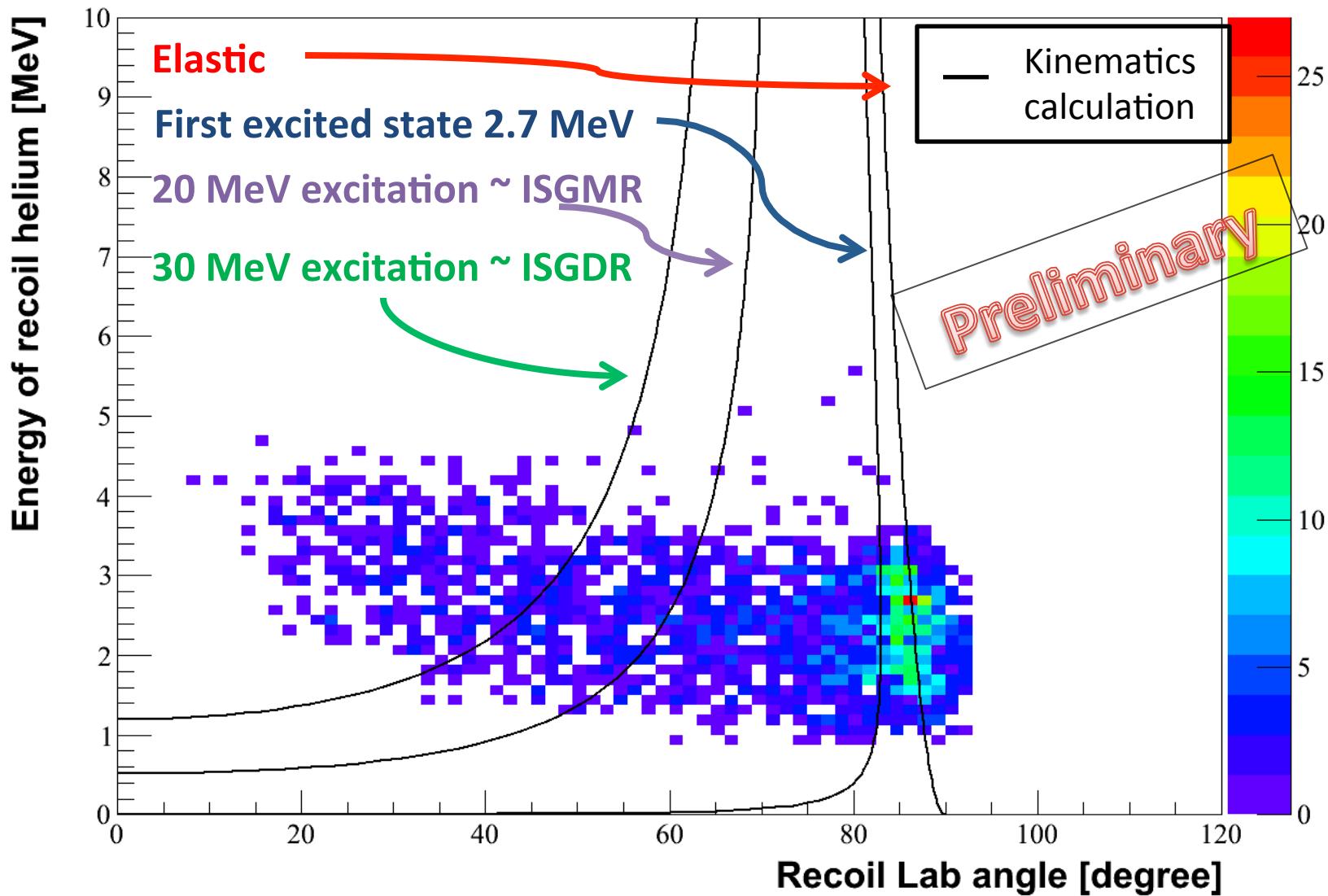
Maximum Amplitude / 2



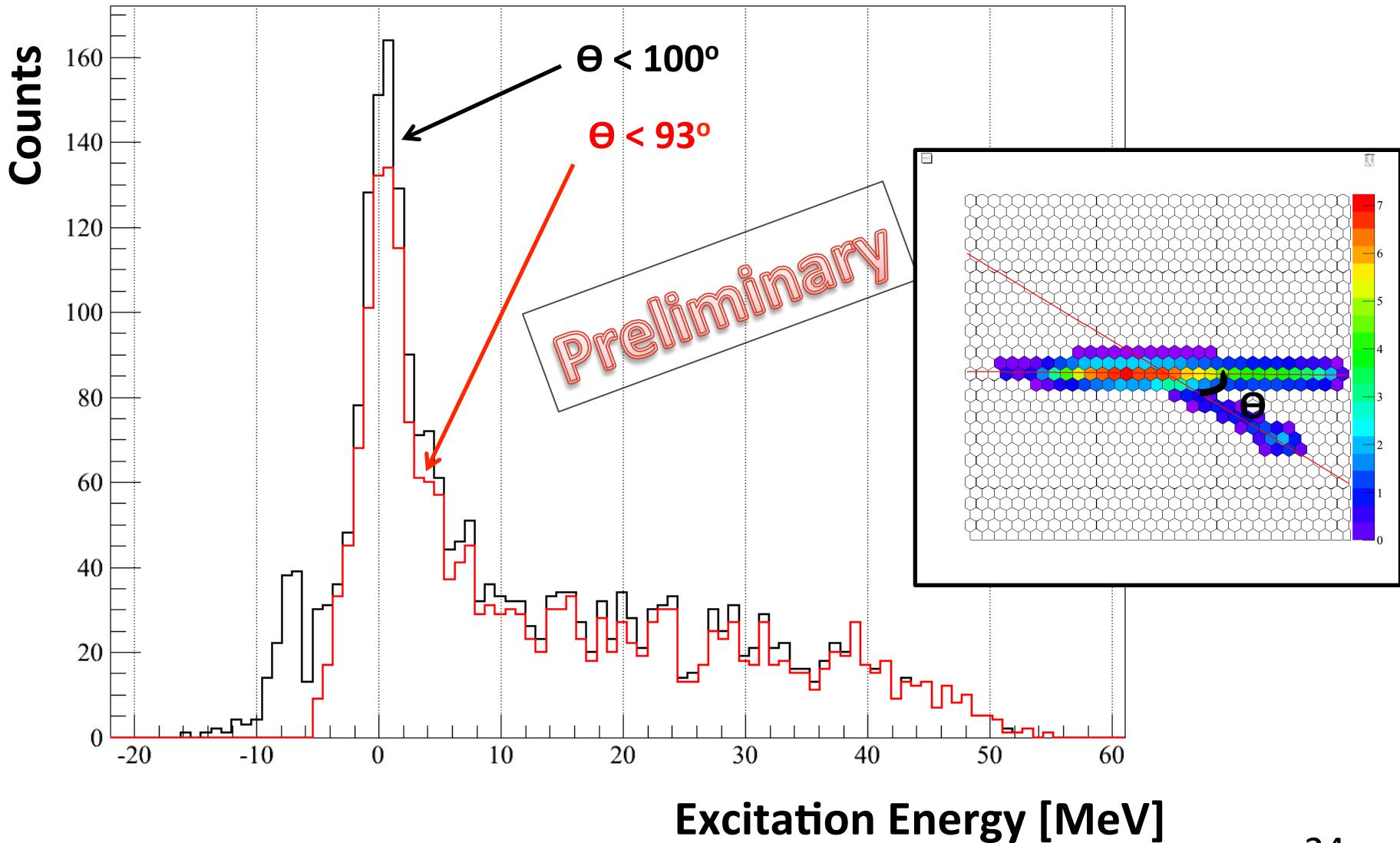
ϕ angle reconstruction



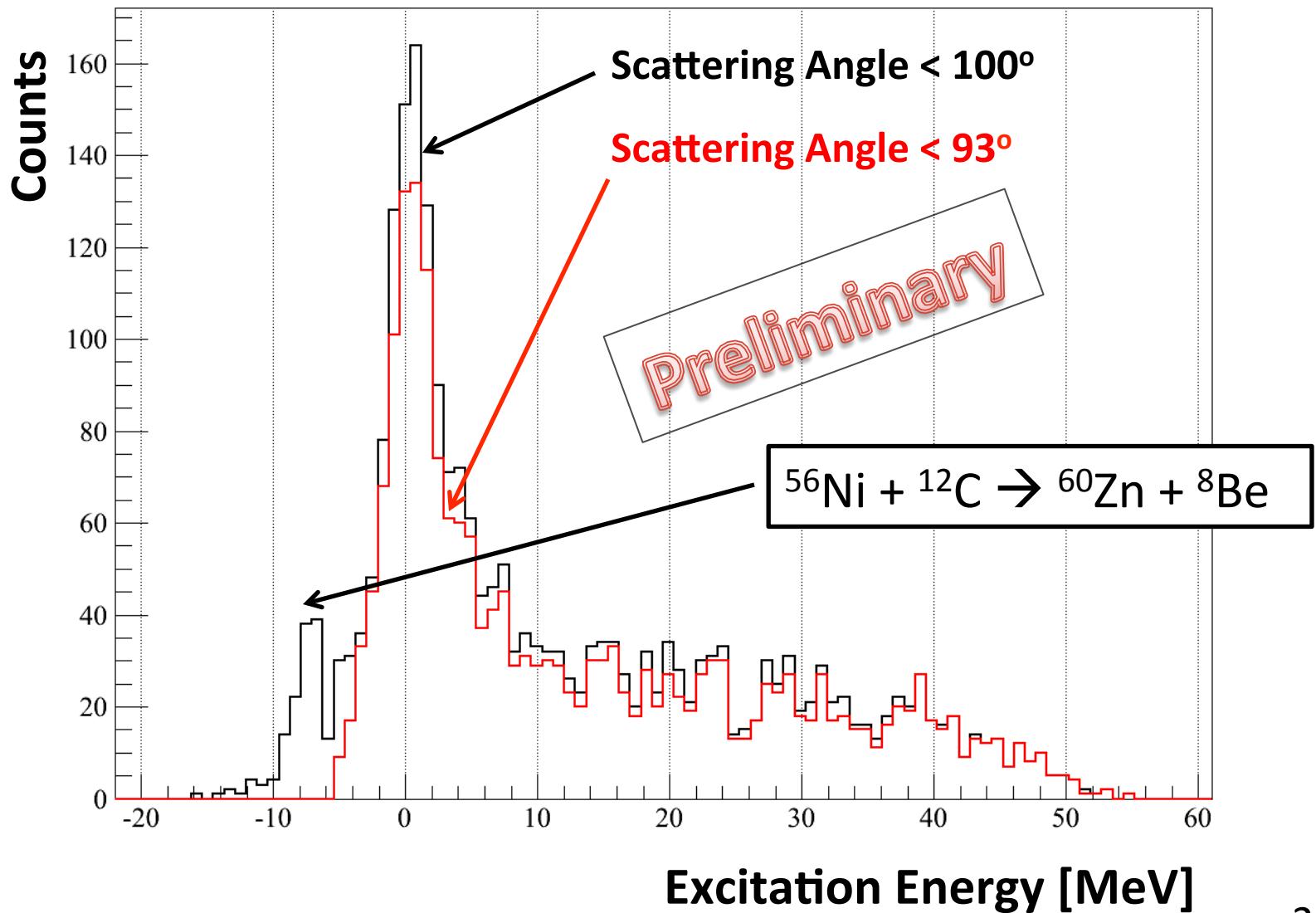
Recoil Helium



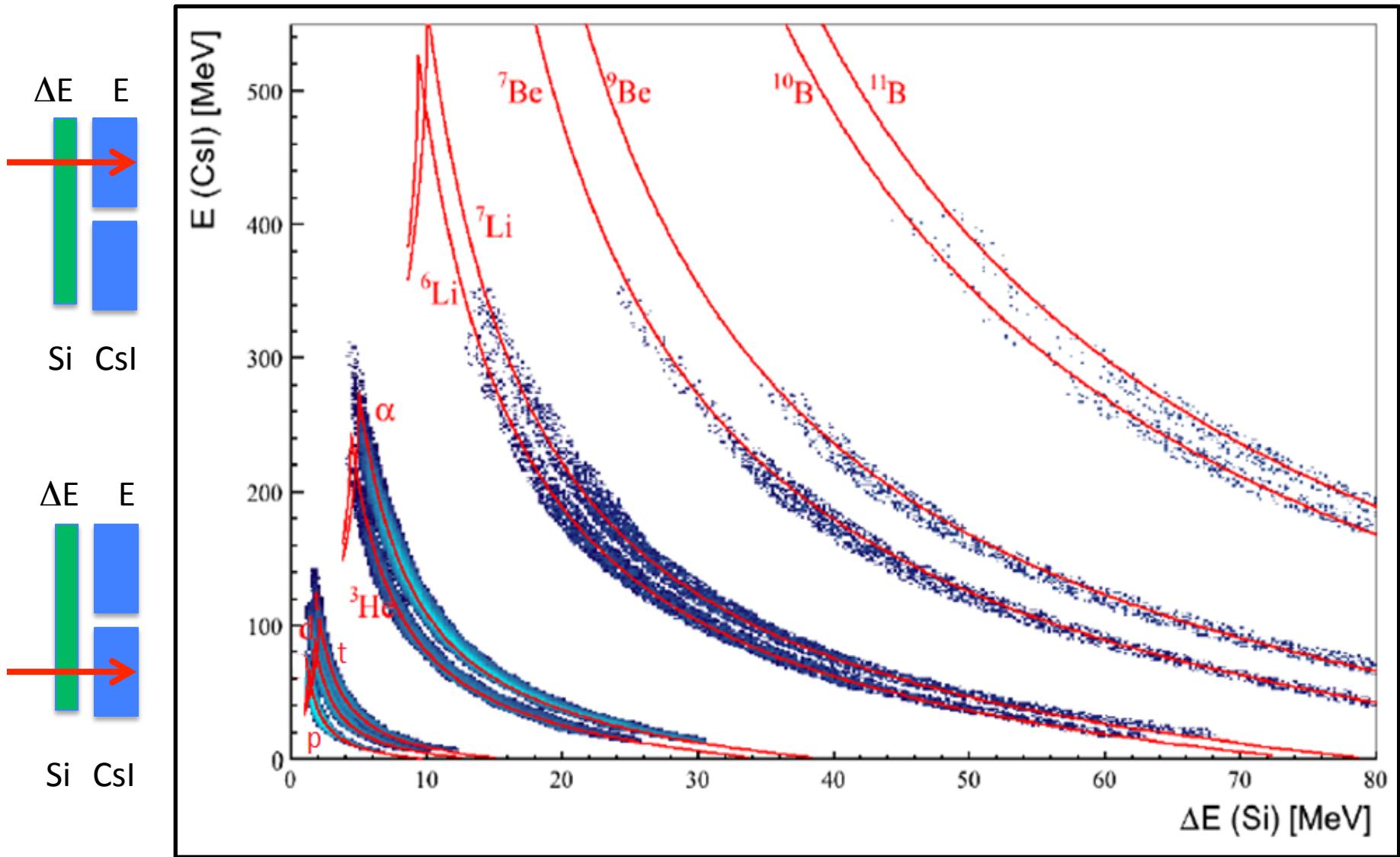
Excitation energy of ^{56}Ni



Excitation energy of ^{56}Ni



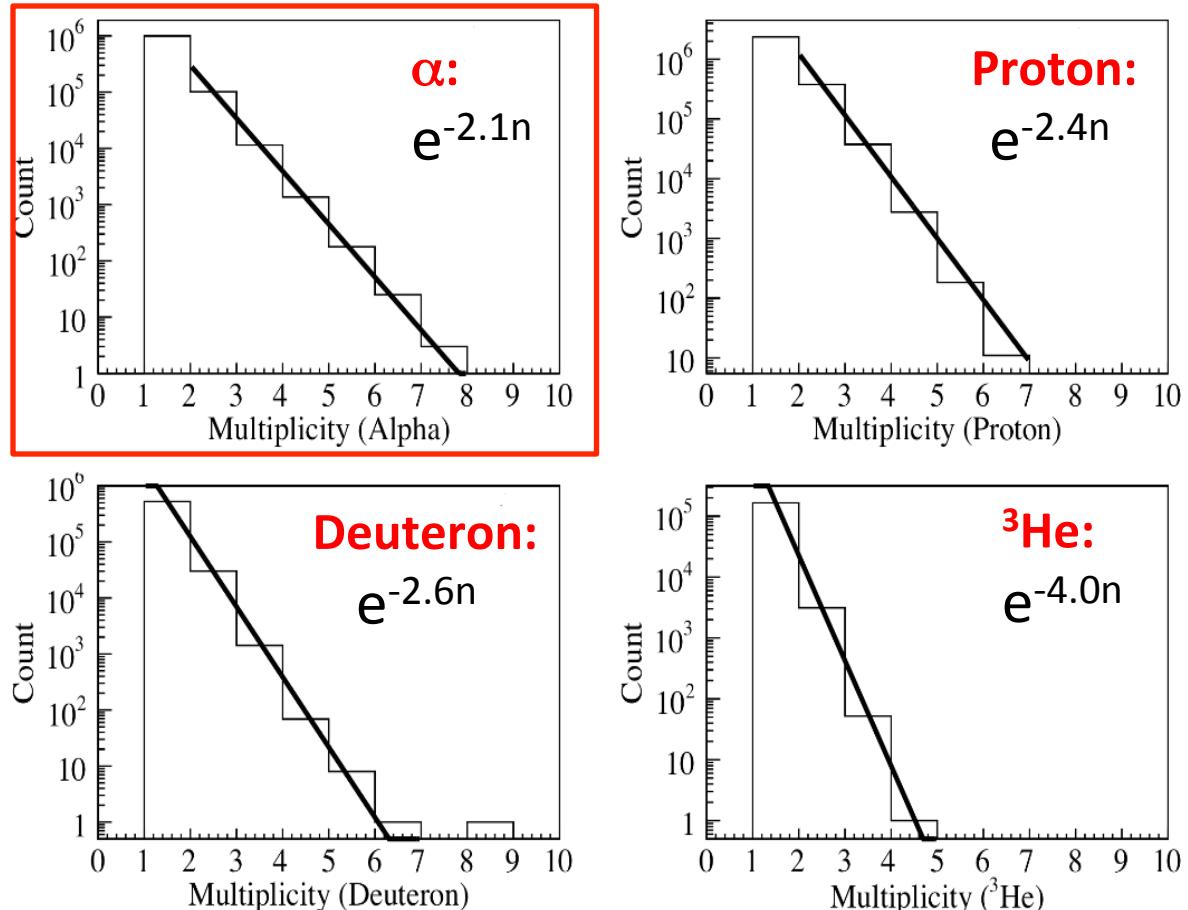
Particle Identification in forward ΔE -E telescope



This analysis is carried out by Prof. H Akimune, Konan University,
Japan

Multiplicity of decay particles

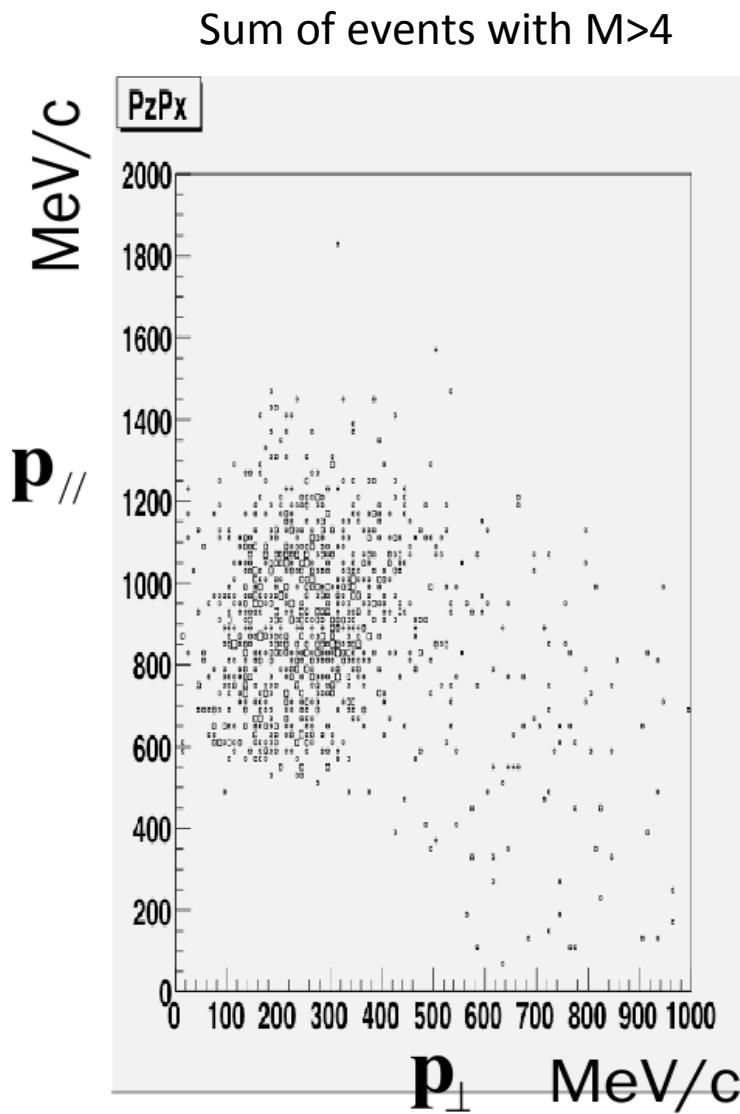
The maximum
multiplicity for α is 7!



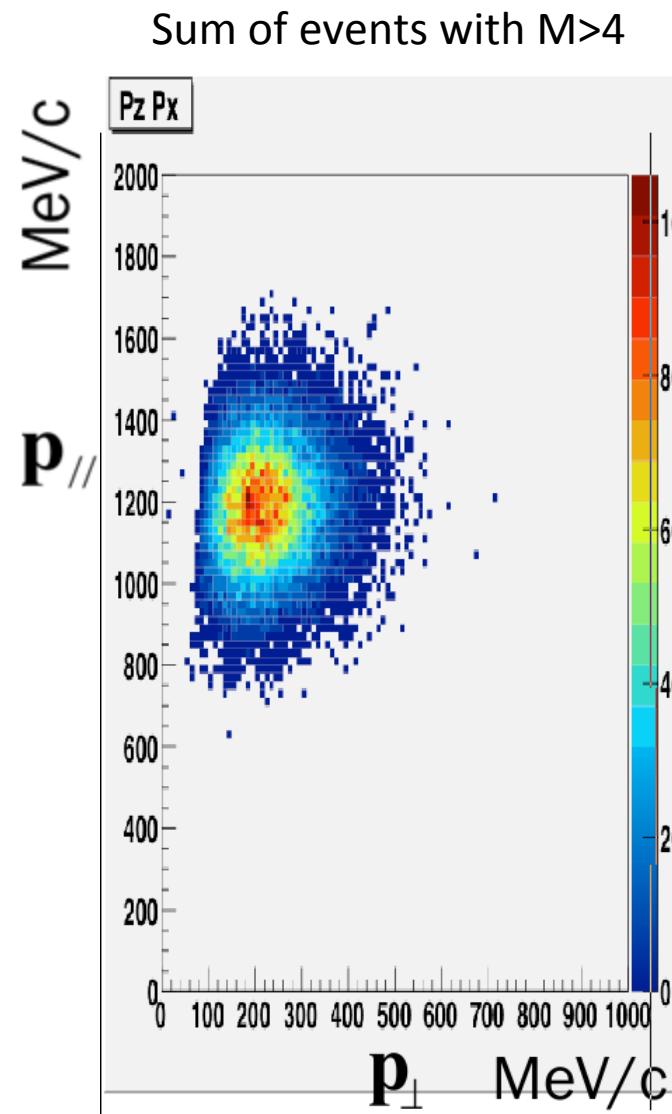
Striking enhancement
of α multiplicity
strongly suggests
existence of α cluster
state in ${}^{56}\text{Ni}$

Momentum distribution of α

Experiment



Simulation for $kT= 5$ MeV, including the efficiency of the detector



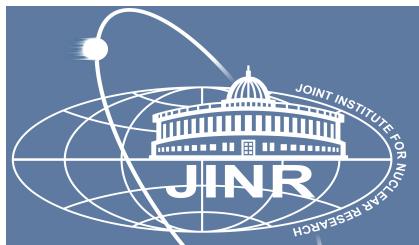
Summary and Outlook

- Compression modes in ^{56}Ni have been studied.
- MAYA active target at GANIL was used.
- Excitation energy of ^{56}Ni (ongoing analysis) is shown.
- Angular distributions of ISGMR and ISGDR of ^{56}Ni have to be obtained.
- Few results of α cluster state in ^{56}Ni (ongoing analysis) are shown.



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lpc
caen

USC
UNIVERSIDADE
DE SANTIAGO
DE COMPOSTELA



 **UNIVERSITY OF
NOTRE DAME**

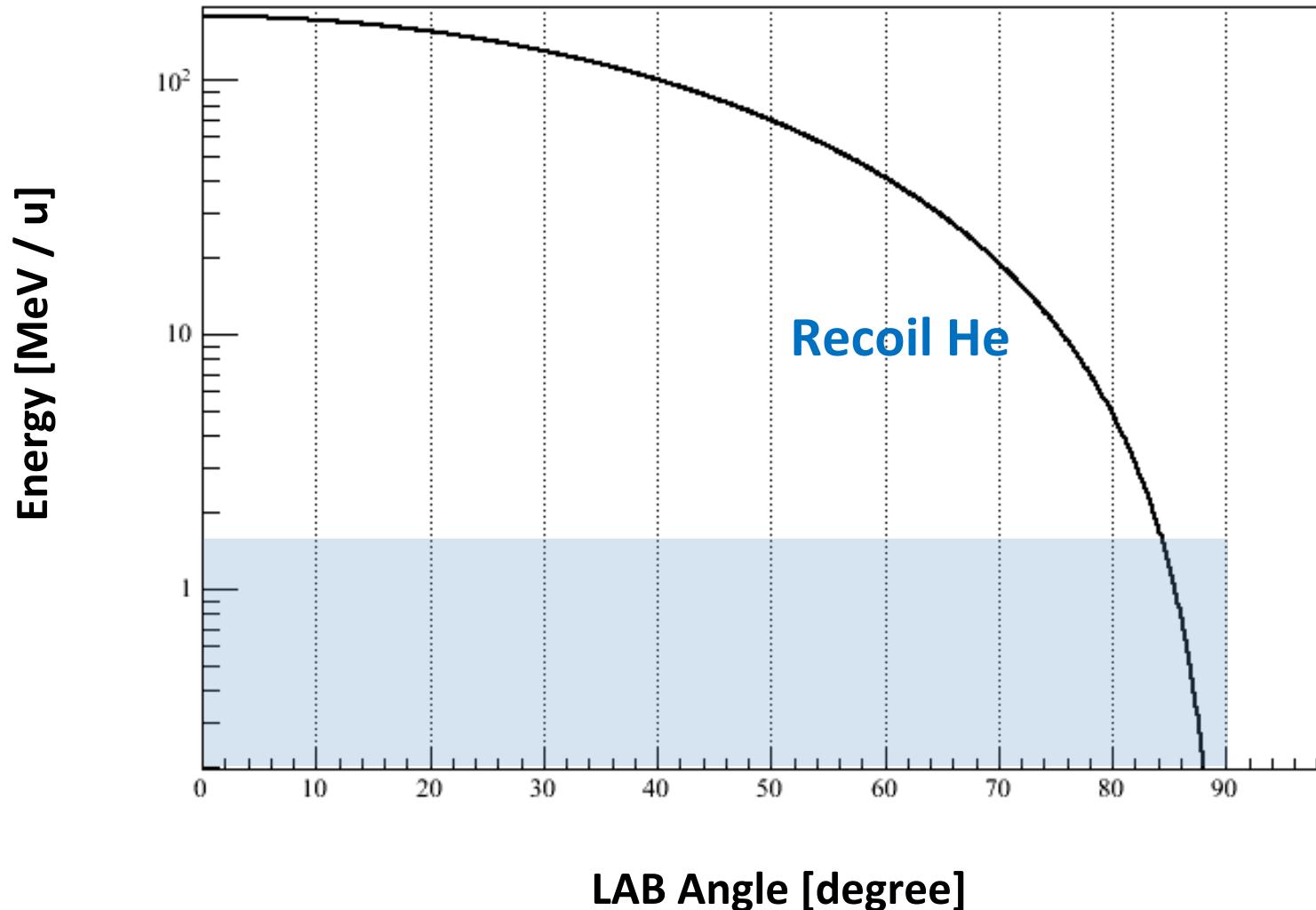
 **OSAKA UNIVERSITY**

 **KVI**

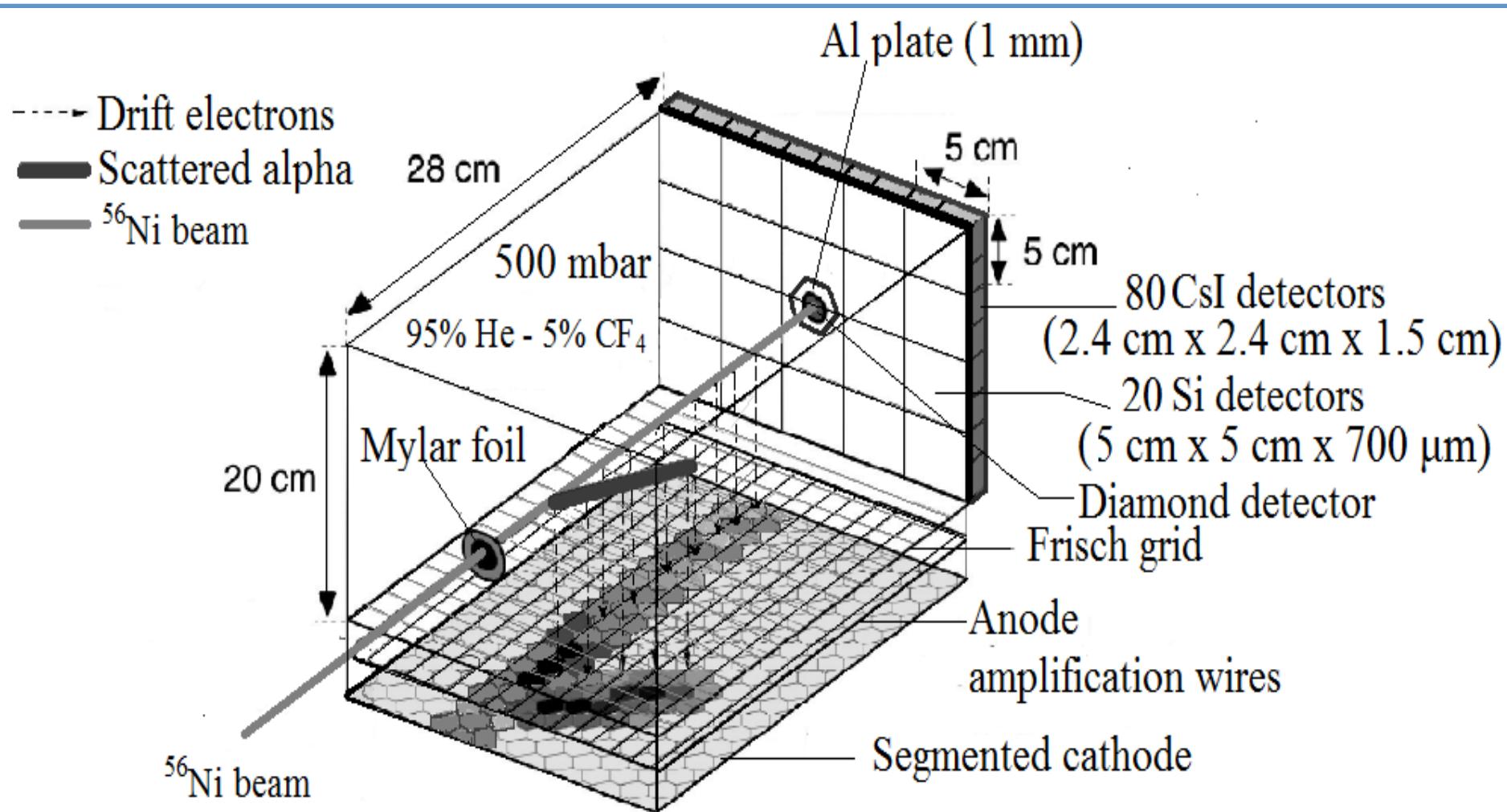
**KATHOLIEKE UNIVERSITEIT
LEUVEN**

 **IPN**
INSTITUT DE PHYSIQUE NUCLÉAIRE
ORSAY

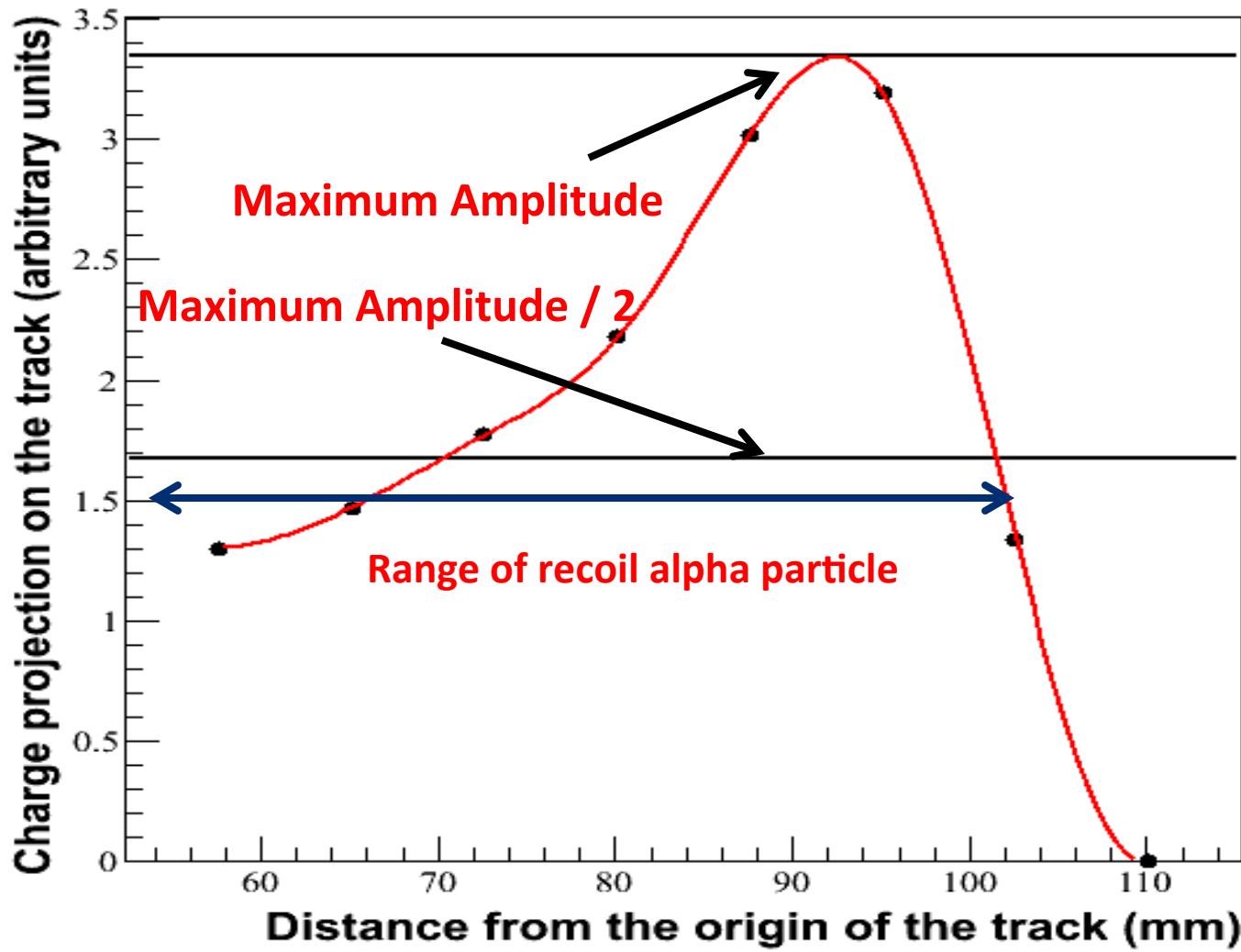
Kinematics curve for ^{56}Ni (α, α) ^{56}Ni with ^{56}Ni at 50 MeV / u



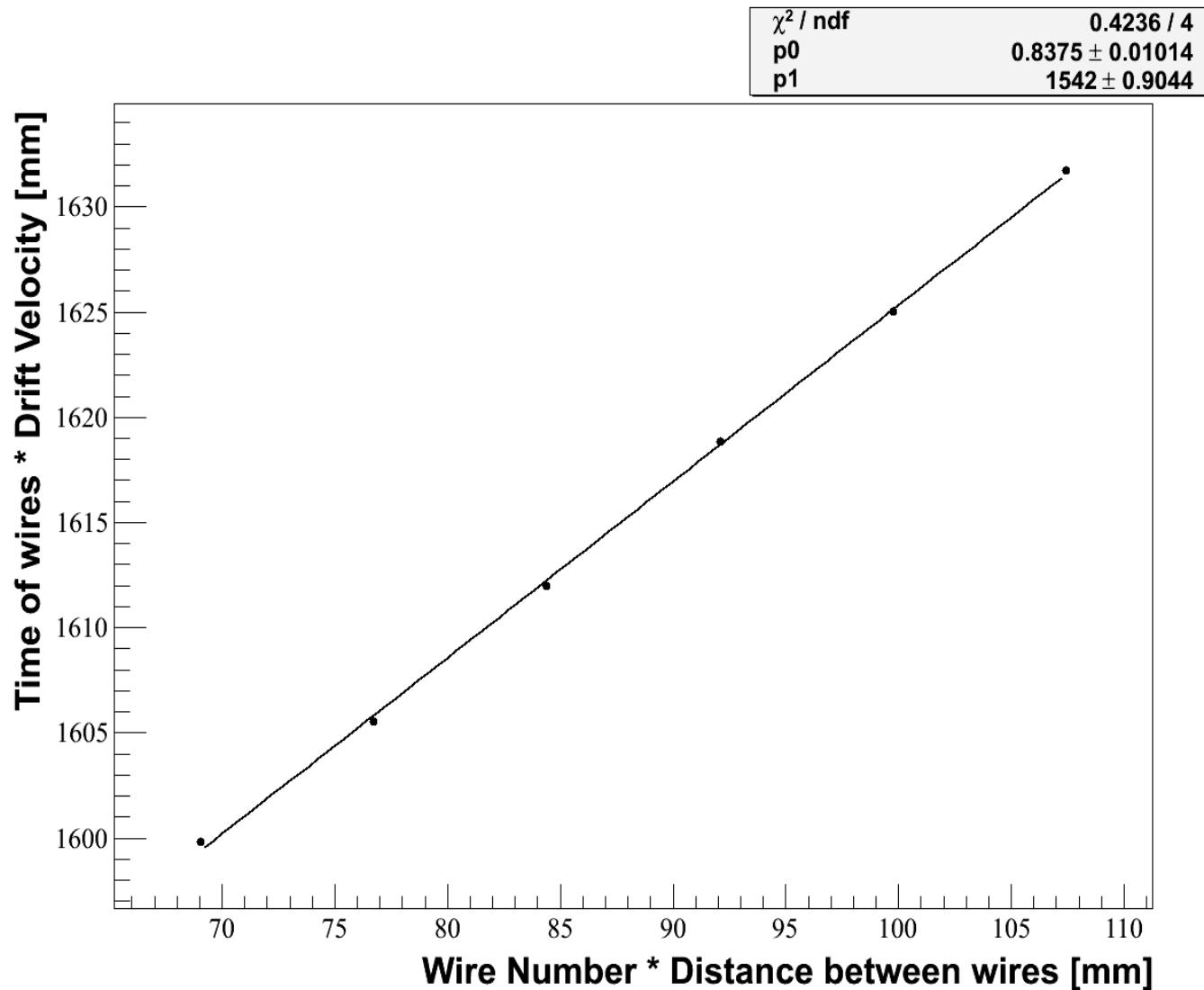
MAYA setup



Range extraction of recoil track

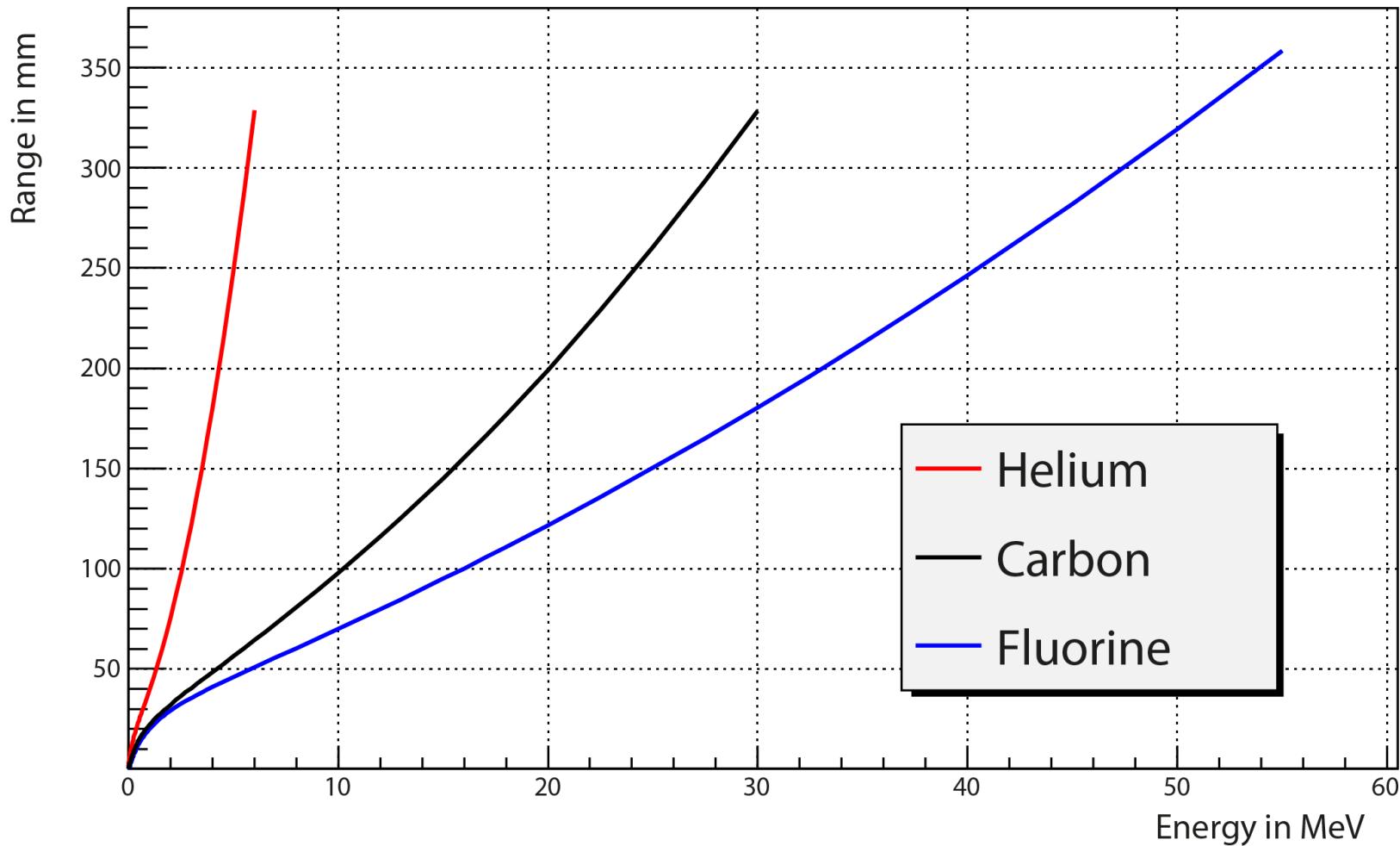


Phi angle reconstruction



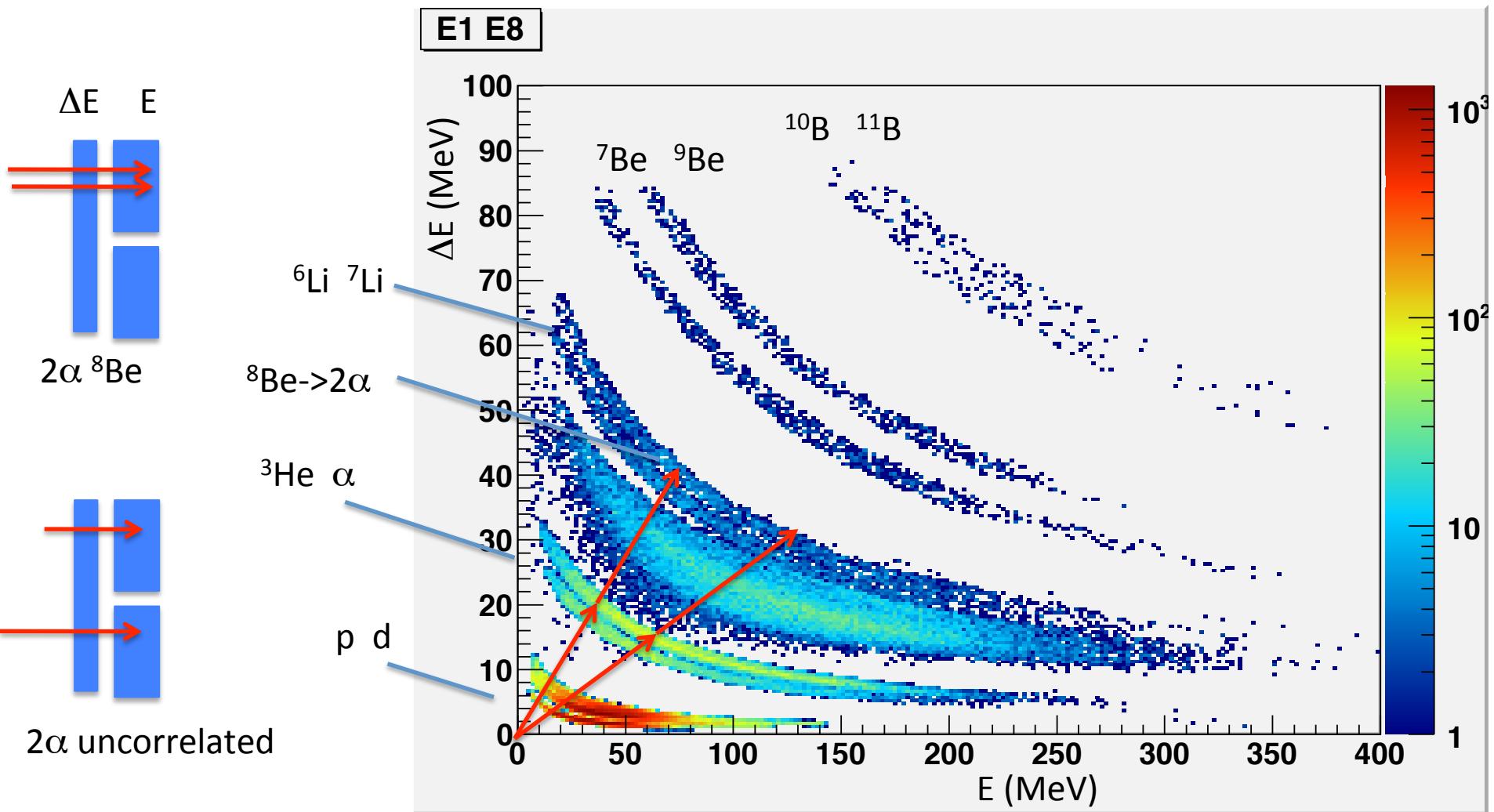
Background Subtraction-SRIM Calculation

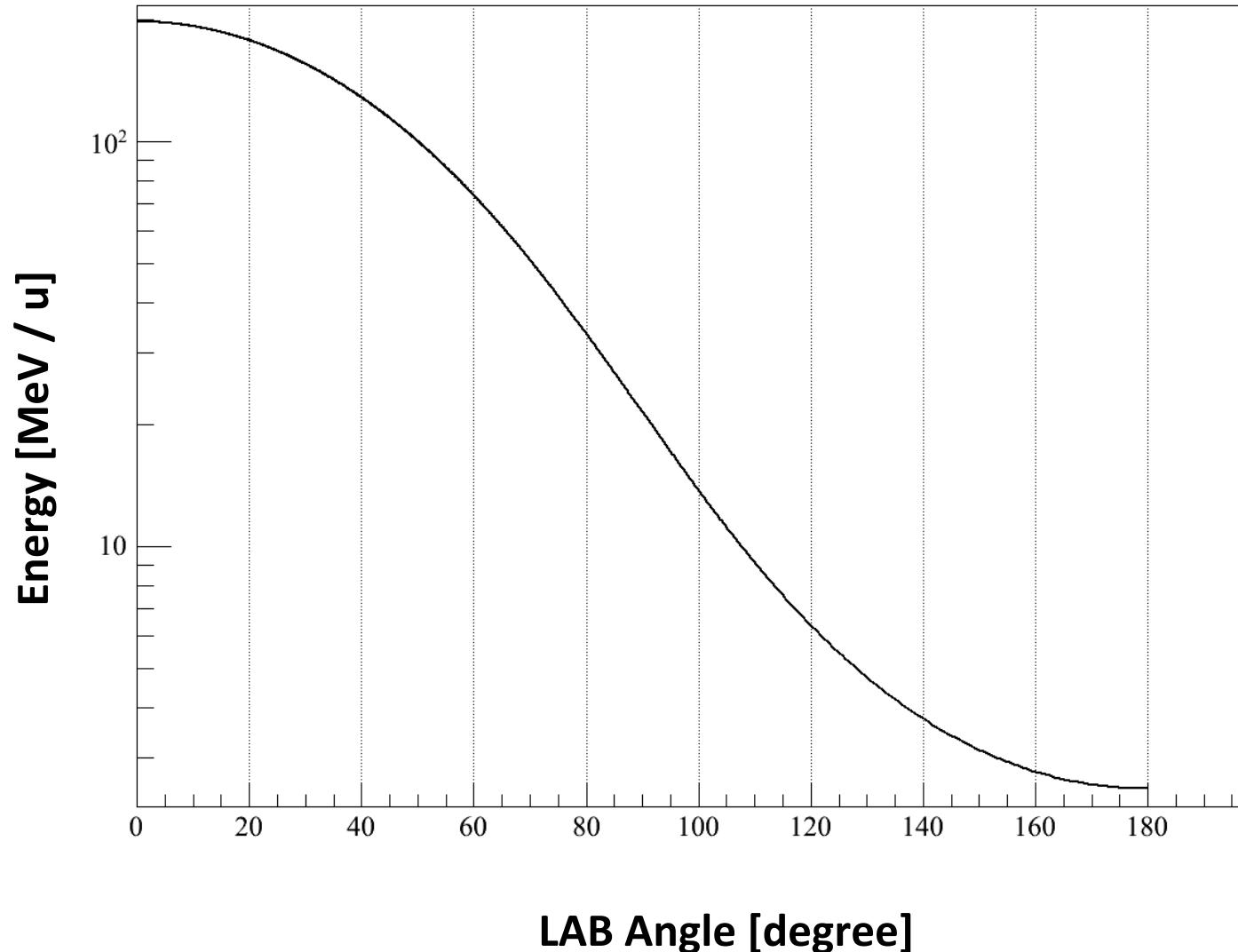
Range of the components of gases in MAYA as a function of energy



Particle identification

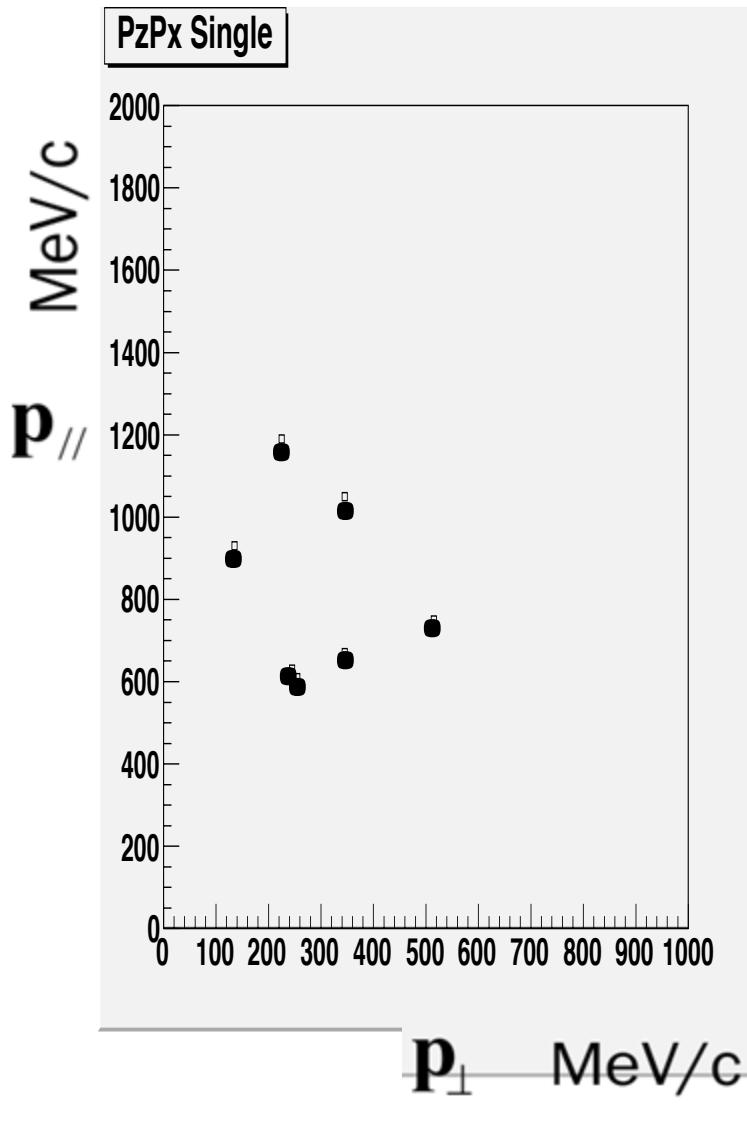
- number of p is the largest,
- however, as for the multiplicity, number of multi α is very large



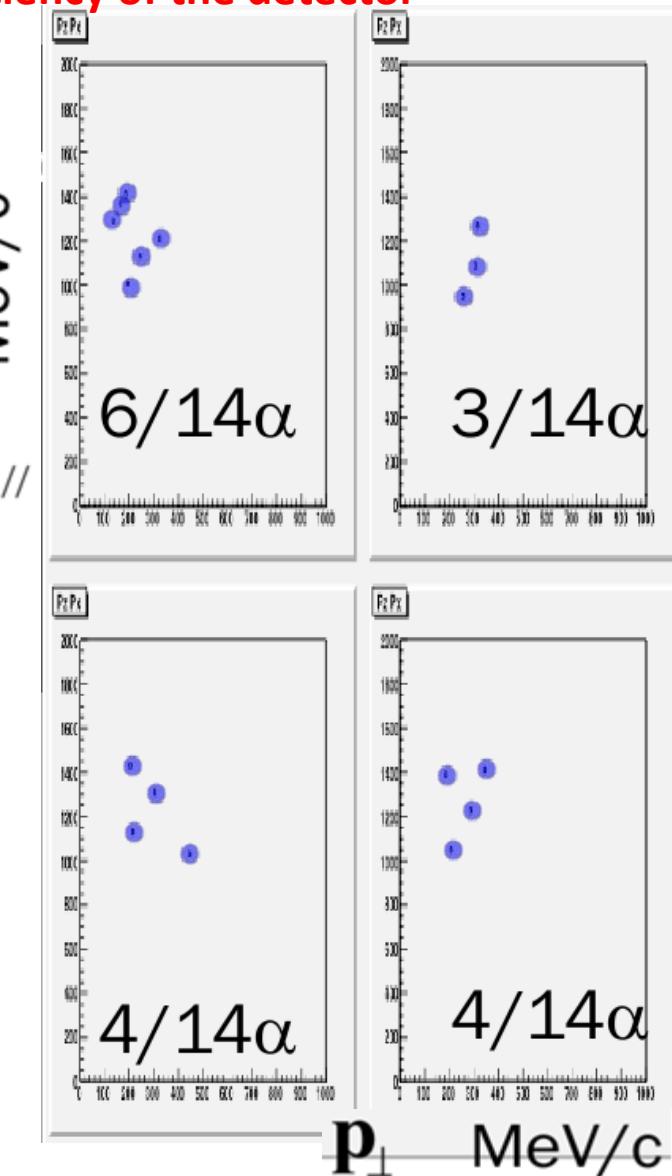


Momentum distribution of α

Experiment



Simulation for $kT = 5$ MeV, including the efficiency of the detector



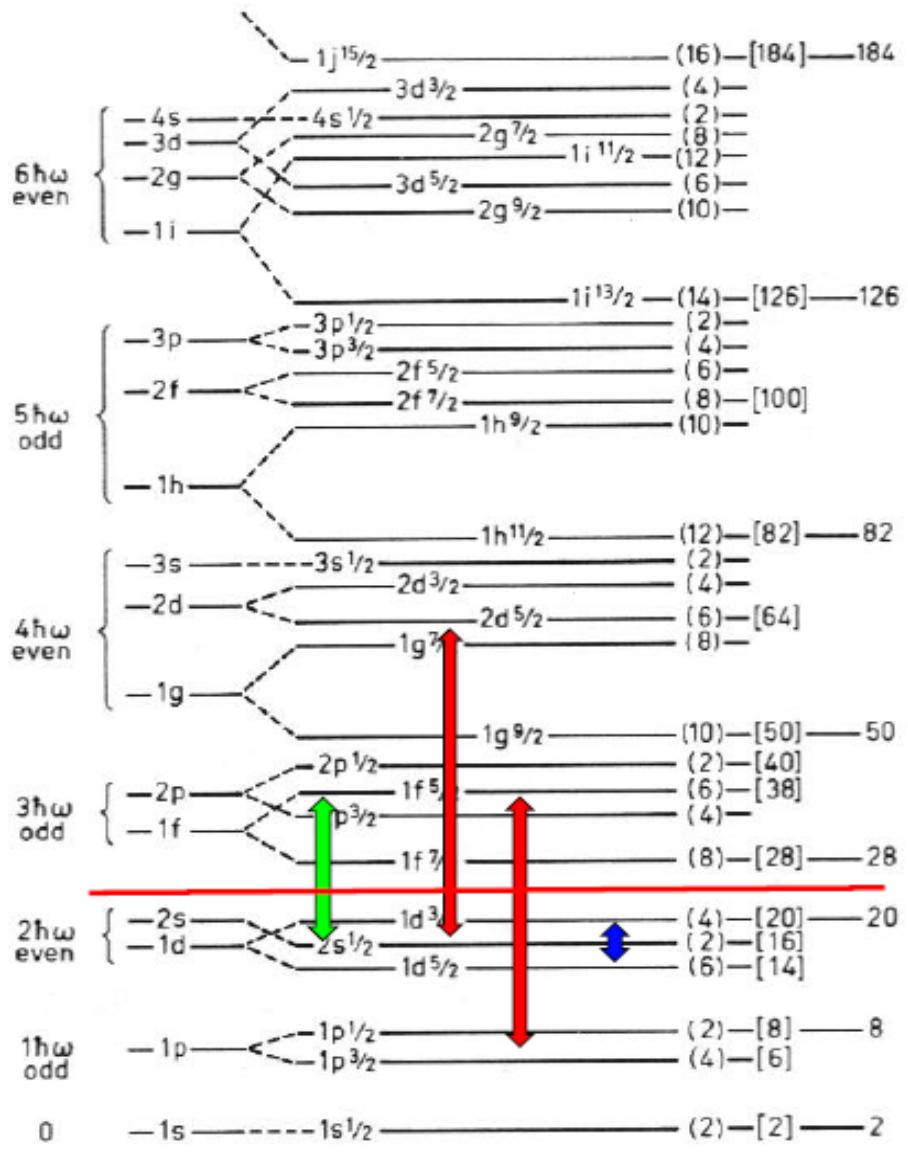
$$K_A = \left[r^2 (d^2(E/A)/dr^2) \right]_{r=R_0}$$

J.P. Blaizot, Phys. Rep. 64 (1980) 171

ISGMR, ISGDR \Rightarrow Incompressibility,
symmetry energy

$$K_A = K_{vol} + K_{surf} A^{-1/3} + K_{sym} ((N-Z)/A)^2 + K_{Coul} Z^2 A^{-4/3}$$

↔ $\Delta N = 1$ E1 (IVGDR)
 ↔ $\Delta N = 2$] E2 (ISGQR)
 ↔ $\Delta N = 0$] E0 (ISGMR)



Decay of giant resonances

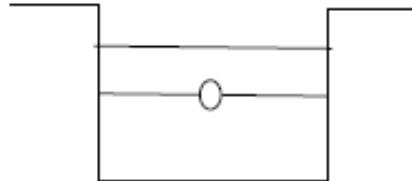
- Width of resonance

$\Gamma, \Gamma^\uparrow, \Gamma^\downarrow (\Gamma^{\downarrow\uparrow}, \Gamma^{\downarrow\downarrow})$

- Γ^\uparrow : direct or escape width

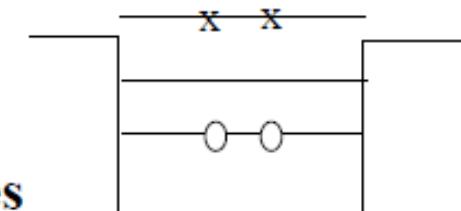
- Γ^\downarrow : spreading width

$\Gamma^{\downarrow\uparrow}$: pre-equilibrium, $\Gamma^{\downarrow\downarrow}$: compound



Γ^\uparrow

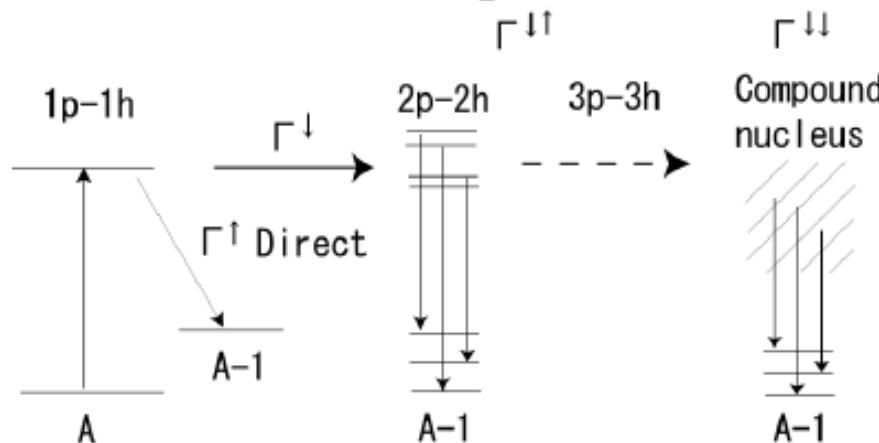
Γ^\downarrow



- Decay measurements

\Rightarrow Direct reflection of damping processes

Allows detailed comparison with theoretical calculations

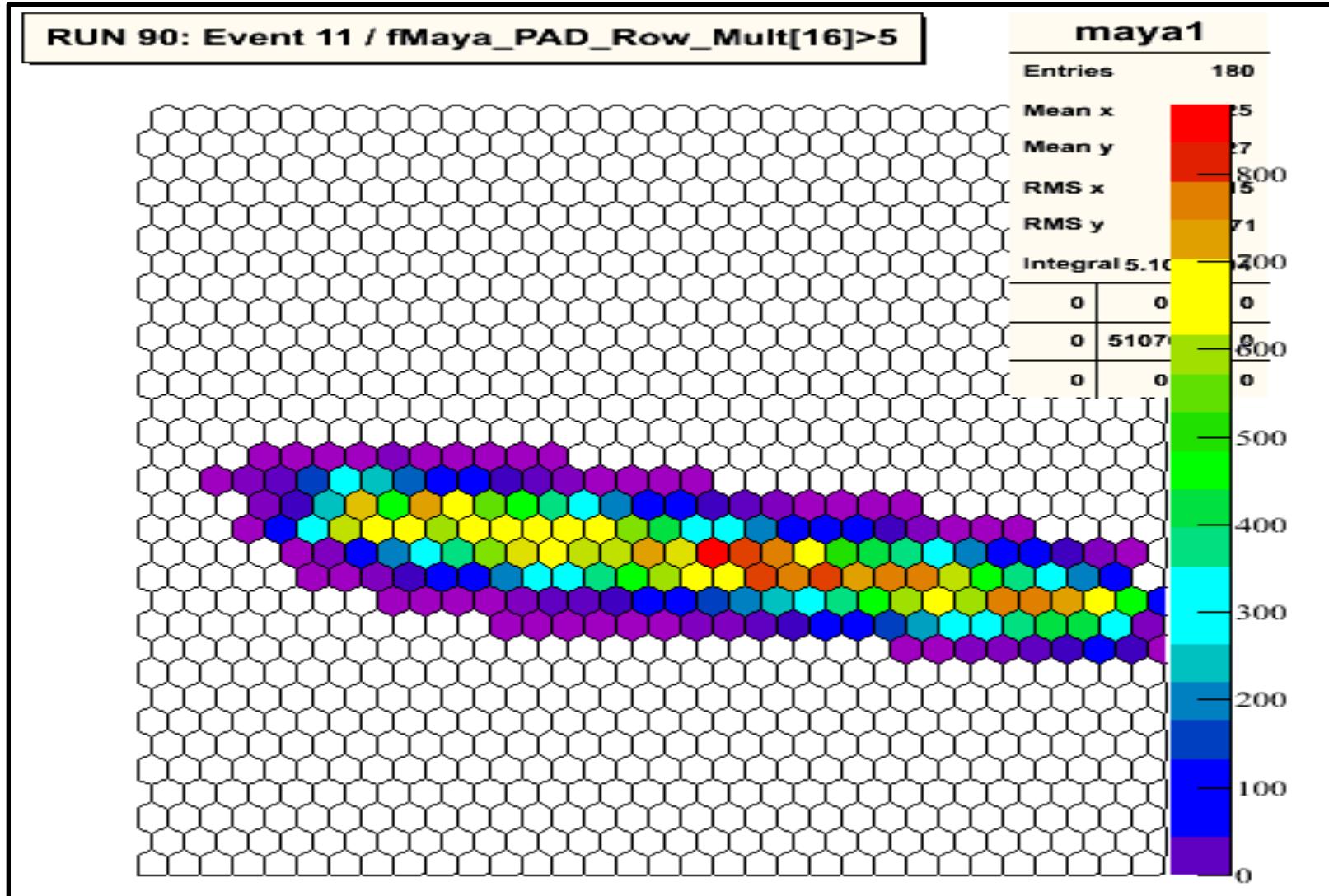


Microscopic picture:

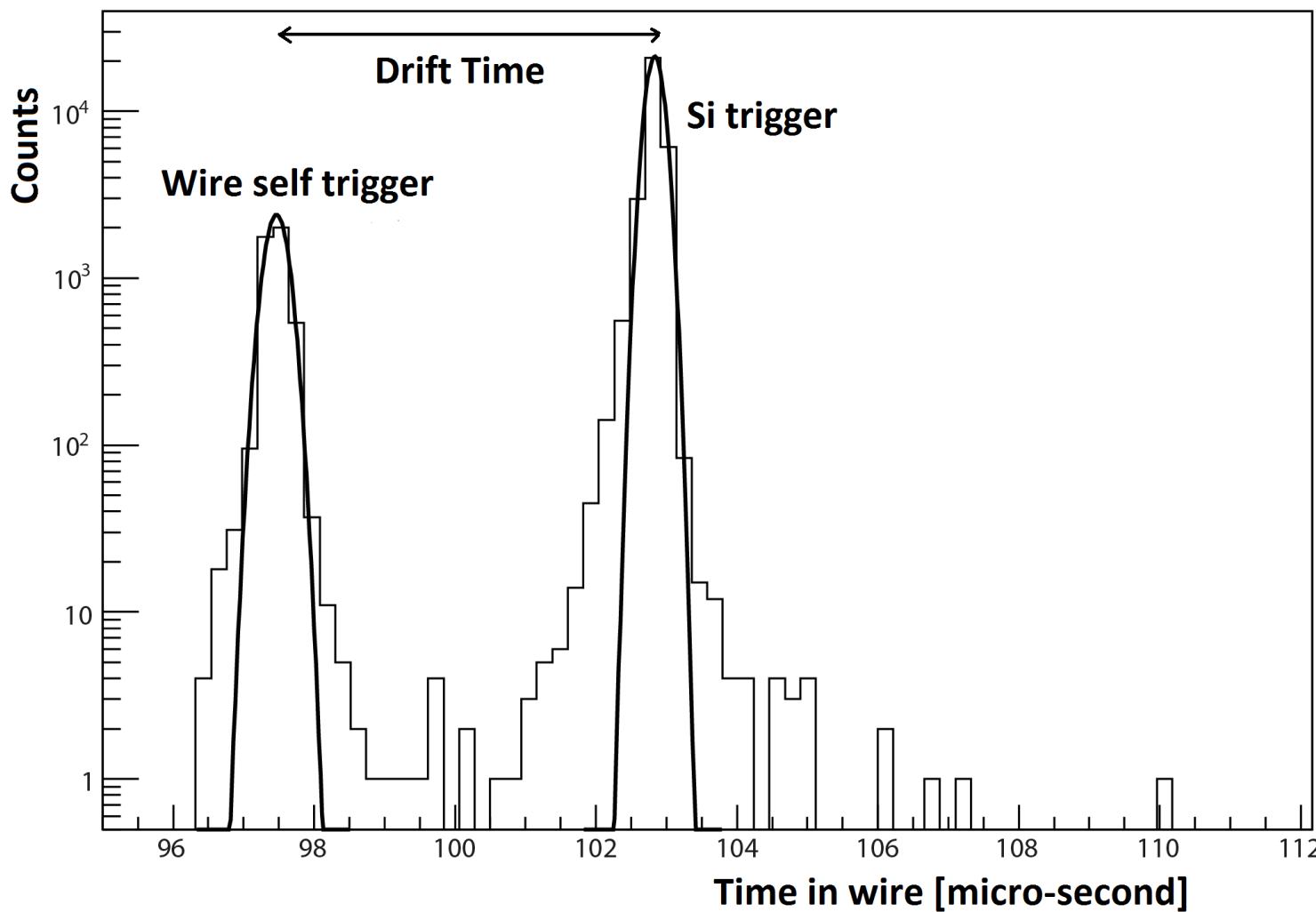
GRs are coherent superposition of 1p-1h excitations induced by the single particle operators

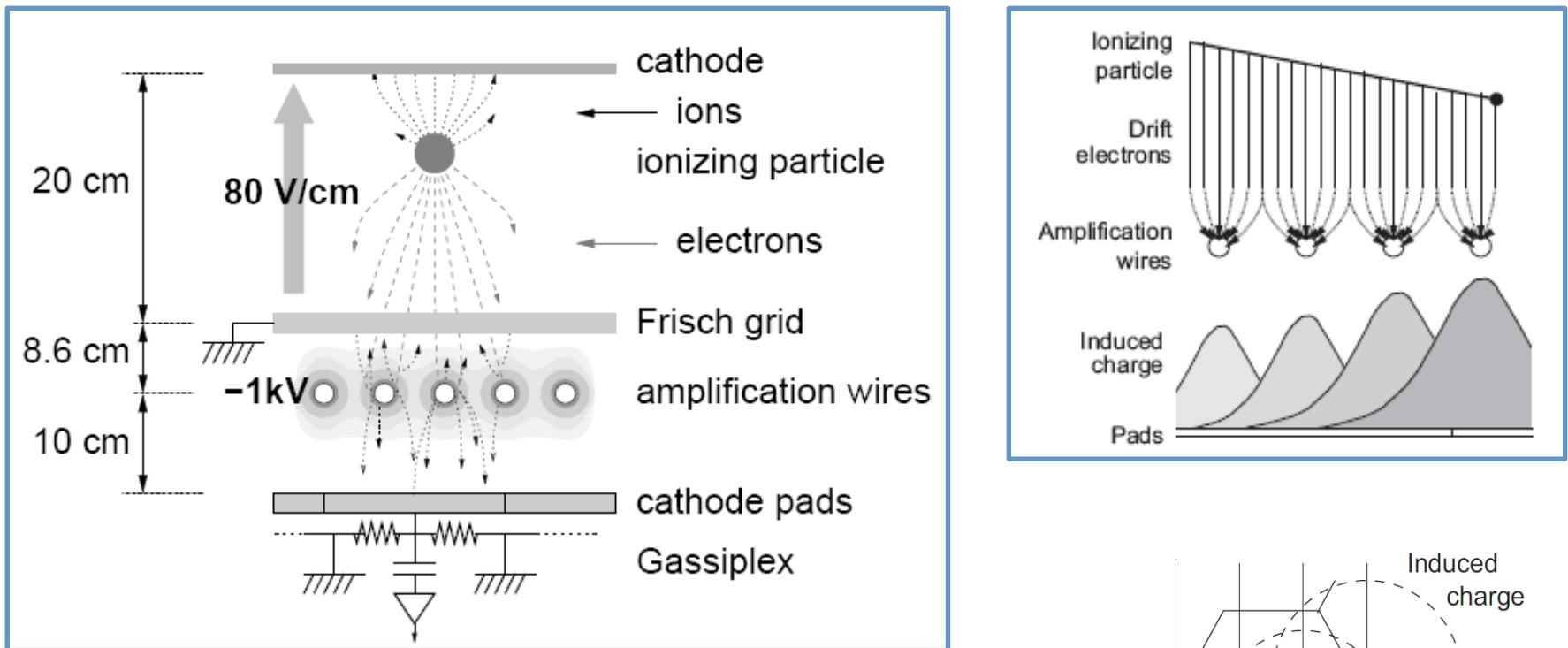
- Excitation energy depends on:
 - Multipole L ($L\hbar\omega$, since the radial operator $\propto r^L$; except for ISGMR and ISGDR, $2\hbar\omega$ and $3\hbar\omega$ respectively)
 - Strength of effective interaction
 - Collectivity
- Exhausts appreciable % of EWSR
- Acquires a width due to coupling to continuum and to the underlying 2p-2h.... configurations

Drift time measurement

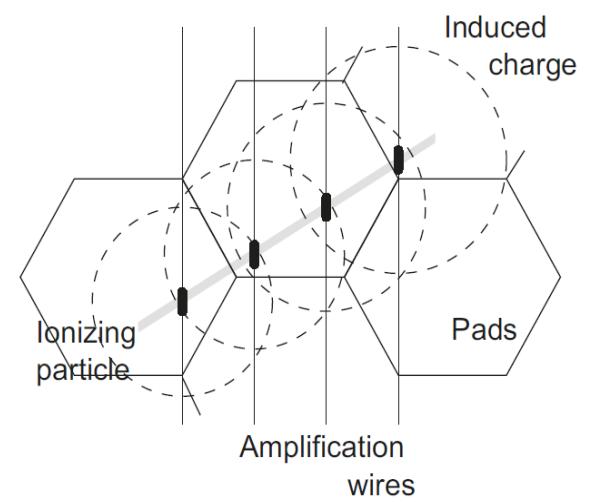


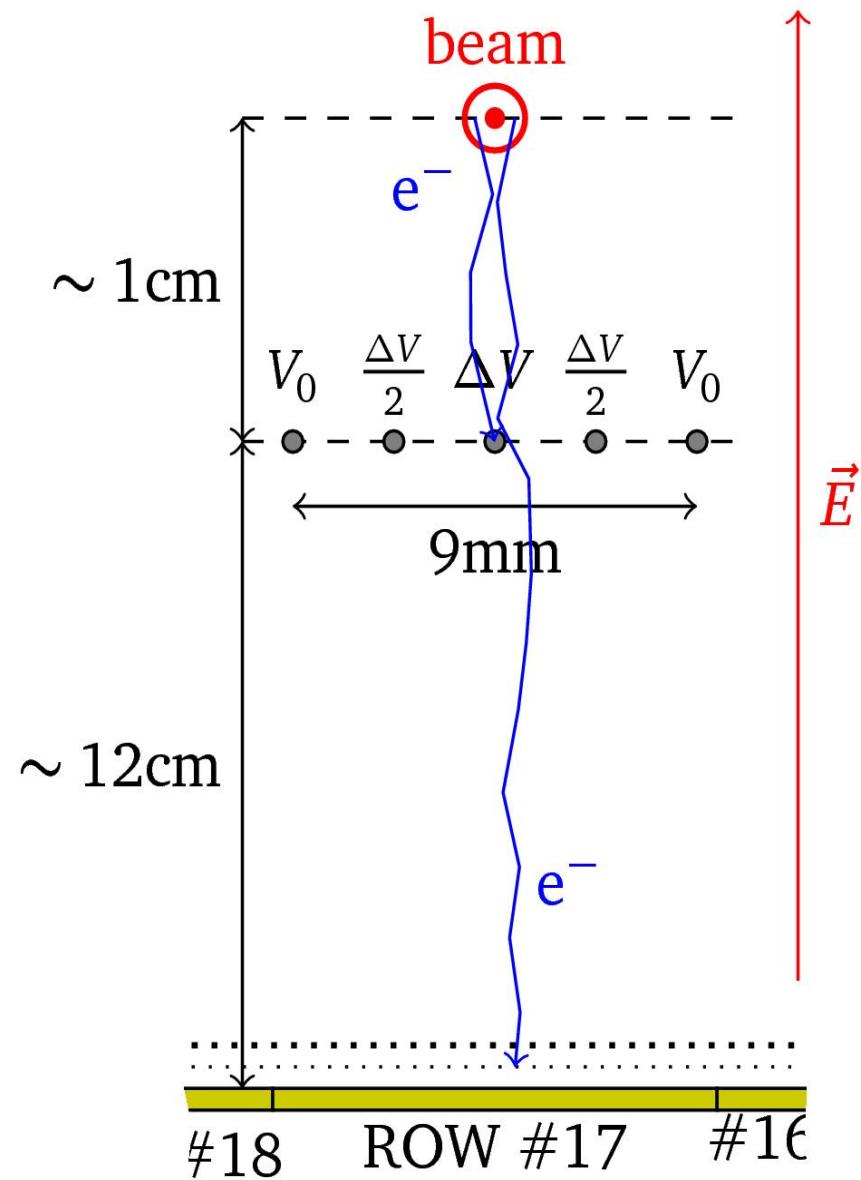
Drift time measurement



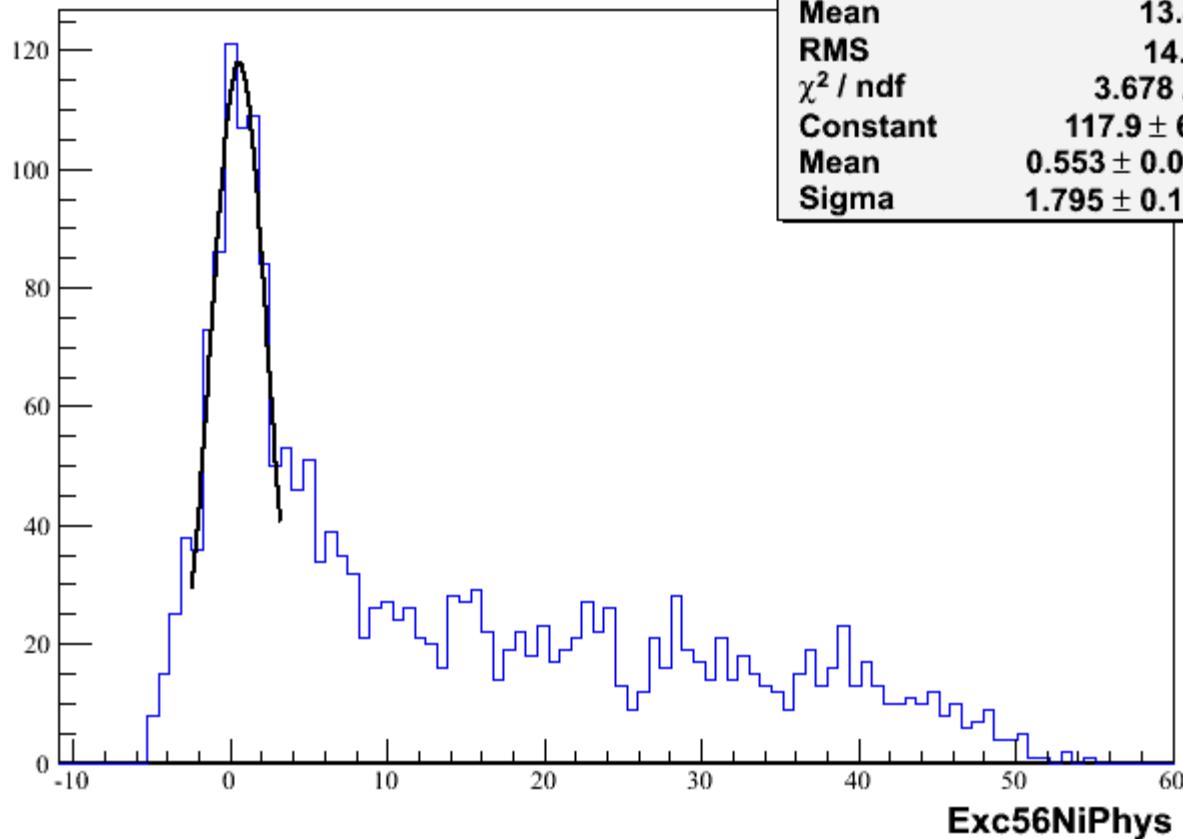


$$\sigma(x, y) = \frac{-Q}{2\pi} \sum_{n=0}^{\infty} \frac{(2n+1)L}{[(2n+1)^2 L^2 + x^2 + y^2]^{3/2}}$$





Exc56NiPhys



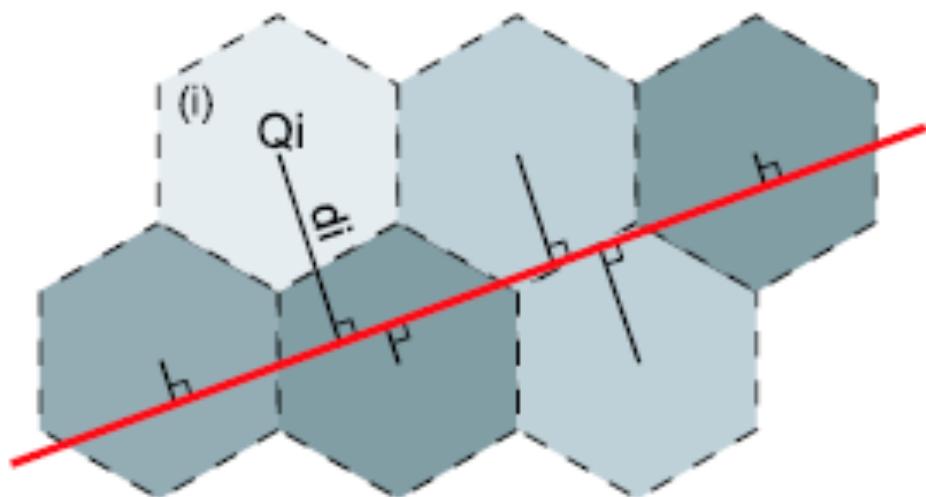


Fig. 9. Principle of the *Global Fitting method*. The orthogonal distance of the center of the pads to the straight line, weighted by their charge is minimized.

$$\chi^2 = \sum_{n=0}^N Q_n \frac{(a_0 x_n + a_1 - y_n)^2}{a_0^2 + 1}$$

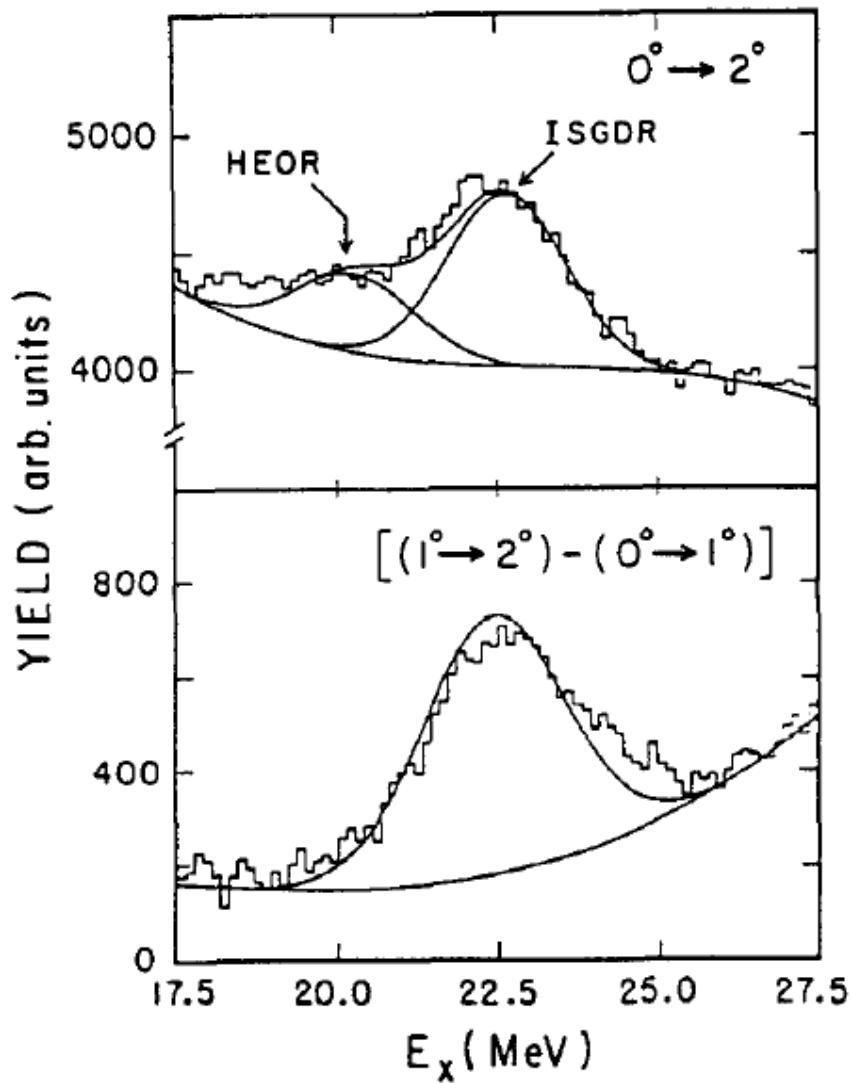


Figure 1. 200 MeV inelastic α -scattering spectra for ^{208}Pb for $(0 \rightarrow 2)^\circ$ (upper panel) and the “difference” spectrum (lower panel). Simultaneous 2-peak + polynomial-background fits to the two spectra are shown superimposed.